



The experimental setup of widely-tunable picosecond pulse generator is shown in Fig.1. The pump source is a tunable laser source (TLS), which is tuned from 1532 nm to 1549 nm to achieve the widest tuning range. The CW output of the TLS is intensity-modulated by a 10 MHz electrical pulse with duty ratio of 1/1000 to produce optical pump pulse with pulsewidth of 100 ps and repetition rate of 10 MHz. Polarization controller PC1 aligns the pump's state of polarization (SOP) with the transmission axis of the MZM. The pump is then amplified by EDFA1, filtered by a 1-nm bandwidth TBPf and further amplified by EDFA2 to produce a high-power, low-noise pump. It then passes through a circulator and the reflected power is observed using a power meter. As measured from the power meter, the stimulated Brillouin scattering (SBS) is low and therefore the phase modulator used in [8] is not required to suppress the SBS. It is then coupled into the cavity for parametric amplification through a wavelength-division multiplexing coupler (WDMC1), which has a cutoff wavelength of 1530 nm. The peak power of the pump is measured to be 25 W at the fiber input. The gain medium inside the cavity is a spool of 40-m dispersion-shifted fiber (DSF), which has nonlinear coefficient of  $2 \text{ W}^{-1}\text{km}^{-1}$ , zero-dispersion wavelength (ZDW) of 1548.7 nm, and dispersion slope of  $0.0693 \text{ ps/nm}^2/\text{km}$ . A 50/50 coupler in the cavity provided 50% feedback and 50% output. The feedback branch is filtered by the WDMC2 with a cutoff wavelength of 1530 nm, so that only the signal (shorter wavelength component) returns to the DSF through WDMC1, and the idler (longer wavelength component) is blocked. As a result, the FOPO is only singly resonant with the signal. The PC2 in the cavity is used to align the SOP of the signal with that of the pump. The total length of the discrete standard single mode fiber (SMF) used to connect the components in the cavity is 20 m, so the total cavity length is 60 m.

There are two conditions need to be satisfied for stable oscillations. One is that the signal pulse should synchronize with the pump pulse, so we need to use an ODL inside the cavity to adjust the cavity length thus the round trip time of the signal matches the pump repetition rate. Since the round trip time of different signal is wavelength dependent because of group-velocity dispersion (GVD), each tuning we need to tune the ODL to synchronize the signal at oscillating wavelength with the pump. The other condition is that the single-pass signal gain should be larger than the cavity loss. The round trip loss of the cavity is measured to be 9.5 dB.

The single-pass signal peak gain of the FOPO in the case of perfect phase matching can be calculated using [7]:

$$G_{dB} \approx 10 \log_{10} \left[ \frac{1}{4} \exp(2\gamma PL) \right] \quad (1)$$

where  $L$  is the fiber length,  $P$  is the peak power of the pump,  $\gamma$  is the nonlinear coefficient. The peak power corresponding to a 9.5-dB gain is calculated to be 22.3 W. The experimental threshold pump peak power is 25 W, matches reasonably well with the calculated value. After the above two conditions are satisfied, stable oscillations can be obtained. The FOPO output spectrum is monitored by an OSA through a 1/99 coupler. WDMC3 is used to filter out the desired signal or idler. The pulsewidth of the signal (idler) is measured using an autocorrelator.

### 3. Results and Discussion

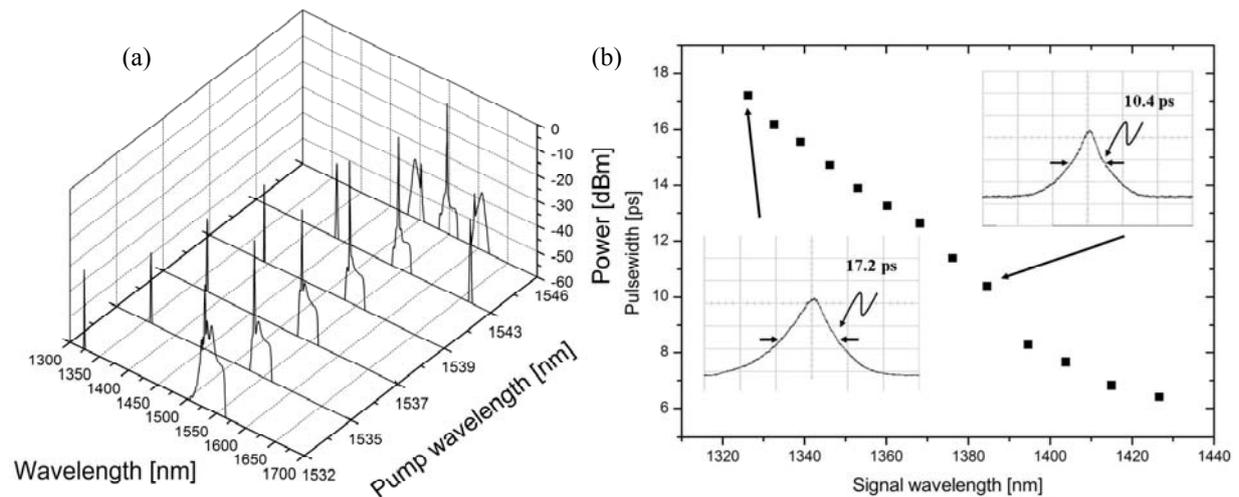


Fig. 2. (a) Optical spectra measured at FOPO output using OSA, (b) Signal pulsewidth versus wavelength. Insets are autocorrelation traces at some wavelengths.

Fig. 2 (a) shows the optical spectra measured at the FOPO output port. Tunability is achieved by tuning the pump wavelength from 1549 nm to 1532 nm, and the pump power is slightly adjusted to maintain almost the same for all

spectrum peaks. As the pump wavelength is decreased, the detuning of the sidebands increases. When the ODL is tuned at the same time to synchronize the signal with the pump, the achievable output tuning range is from 1320 nm to 1700 nm measured by the OSA. Since the operating wavelength of our OSA ends at 1700 nm, the idler beyond this wavelength is not shown in Fig. 2 (a). The longest idler wavelength calculated from the experimental measured shortest signal wavelength is 1820 nm. Therefore, the tuning range of our FOPO is from 1320 nm to 1820 nm, which is as wide as 500 nm. The tuning range is larger than those picosecond FOPOs reported in [9, 11]. There are three factors which may combine to limit the further enhancement of the tuning range. Firstly, the DSF is not polarization maintained, and random birefringence may decrease the four-wave mixing (FWM) efficiency at large detuning [12]; secondly, the fluctuations of the ZDW along the fiber will also limit the tuning range of this kind of FOPO [13]; the third factor is the walk-off between the signal and pump, which increases with their wavelength separation and thereby decreases their interaction length.

Fig. 2 (b) shows the output pulsewidth as a function of signal wavelength. Insets are autocorrelation traces at some wavelengths measured by the autocorrelator. The real full-width half-maximum (FWHM) pulsewidth is calculated by assuming a  $\text{sech}^2$  pulse shape, multiplied the FWHM correlation width using a deconvolution factor of 0.648. The pulsewidth increases from 7 ps to 17 ps when the signal detunes further from the pump, thus the walk-off between the signal and the pump becomes larger which broaden the signal pulsewidth. The output pulses are narrower than the pump pulses (100 ps) because of the pulse narrowing effect of FOPA [4].

#### 4. Conclusion

In conclusion, an all-fiber widely-tunable picosecond optical parametric oscillator based on DSF was demonstrated. The output was continuously tunable from 1320 nm to 1820 nm, when the pump wavelength was tuned between 1549 nm to 1532 nm. Pulses were generated with pulsewidth from 7 ps to 17 ps. This scheme has the potential to be a cost-effective source in generation short pulse for ultrafast communication in the non-conventional wavelength bands.

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#### 6. References

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