

Anomalous Light Cones and Valley Optical Selection Rules of Interlayer Excitons in Twisted Heterobilayers

Hongyi Yu,¹ Yong Wang,¹ Qingjun Tong,¹ Xiaodong Xu,^{2,3} and Wang Yao^{1,*}

¹*Department of Physics and Center of Theoretical and Computational Physics, The University of Hong Kong, Hong Kong, China*

²*Department of Physics, University of Washington, Seattle, Washington 98195, USA*

³*Department of Materials Science and Engineering, University of Washington, Seattle, Washington 98195, USA*

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We show that, because of the inevitable twist and lattice mismatch in heterobilayers of transition metal dichalcogenides, interlayer excitons have sixfold degenerate light cones anomalously located at finite velocities on the parabolic energy dispersion. The photon emissions at each light cone are elliptically polarized, with the major axis locked to the direction of exciton velocity, and helicity specified by the valley indices of the electron and the hole. These finite-velocity light cones allow unprecedented possibilities for optically injecting valley polarization and valley current, and the observation of both direct and inverse valley Hall effects, by exciting interlayer excitons. Our findings suggest potential excitonic circuits with valley functionalities, and unique opportunities to study exciton dynamics and condensation phenomena in semiconducting 2D heterostructures.

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Monolayer group-VIB transition metal dichalcogenides (TMDs) are attracting interest for their visible-range direct band gap that promises optoelectronics [1–5], and the exotic valley physics including the valley Hall effect [6,7], the valley magnetic moment [8–11], and the valley optical selection rules [6,12–16]. Because of the strong Coulomb interaction, the optical response is dominated by excitons, the hydrogenlike bound state of an electron-hole pair. The demonstrated electrostatic tunability and optical controllability of valley configurations of excitons in monolayer TMDs have implied new optoelectronic device concepts not possible in other material systems [13–17].

Stacking different TMD monolayers to form van der Waals heterostructures opens up a new realm to extend their already extraordinary properties [18]. MoX₂/WX₂ (X = Se, S) heterobilayers have been realized [19–25] which feature a type-II band alignment with the conduction (valence) band edges located in MoX₂ (WX₂) layer. Exciton then has the lowest energy in an interlayer configuration, from which luminescence is observed [19–21]. Because of the spatially indirect nature, interlayer excitons in MoSe₂/WSe₂ heterobilayers have shown long lifetimes exceeding a nanosecond, repulsive interaction, and electrostatically tunable resonance [19], all of which are highly desirable for the realization of excitonic circuits and condensation [26–29]. An unprecedented aspect of this interlayer exciton system is that the heterobilayers, in general, have incommensurate structures due to lattice mismatch and twist in the stacking which, together with the valley physics inherited from the monolayers, bring in radically new properties.

Here we find that interlayer excitons in twisted MoX₂/WX₂ heterobilayers have unique sixfold degenerate

light cones, at finite center-of-mass velocities on the parabolic energy dispersion. In these light cones, an interlayer exciton can directly interconvert with an elliptically polarized photon, with helicity depending on the electron and hole valley indices. These light cones allow resonant optical injection of excitonic valley polarization and valley current with versatile controllability, as well as the observation of valley Hall and inverse valley Hall effects of interlayer excitons. Contrary to known exciton systems, here the exciton lifetime is long at the energy minimum, while the optical injection and probe are allowed at finite velocities, suggesting unique opportunities for studying exciton dynamics and condensation phenomena. For lattice matching heterobilayers of AA-like (0° twist) or AB-like (60° twist) stacking, all light cones merge into a single one, and the quantum interference leads to dependence of the brightness and the polarization selection rules on the interlayer translation.

The optically active interlayer excitons are the ones with electrons (holes) from the $\pm\mathbf{K}'$ ($\pm\mathbf{K}$) valleys in MoX₂ (WX₂), where layer hybridization is substantially quenched by the large band offsets between the layers [3,30–35]. The term *light cones* here refers to the exciton phase-space regions where interconversion with photon is directly allowed (without phonon or impurity assistance). The interlayer exciton in heterobilayers is distinct from existing systems since a small twist and/or lattice mismatch shifts the electron valleys away from the hole valleys [Fig. 1(a)]. The small mismatch between the electron and hole valleys can nevertheless be compensated for by the exciton's kinematic momentum \mathbf{Q} [Fig. 1(b)], to allow radiative recombination with a direct transition. This happens at $\mathbf{Q} = \pm\mathbf{Q}_0$, $\pm\hat{C}_3\mathbf{Q}_0$, or $\pm\hat{C}_3^2\mathbf{Q}_0$, where $\mathbf{Q}_0 \equiv \tau\mathbf{K} - \tau'\mathbf{K}'$ is

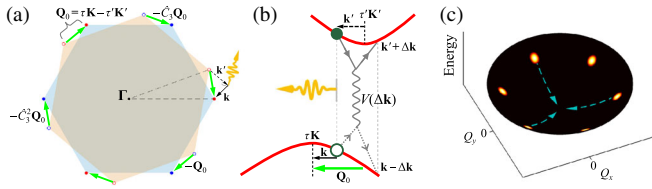


FIG. 1 (color online). (a) With a small twist and/or lattice mismatch, the Brillouin zone corners of the MoX_2 layer ($\tau\mathbf{K}'$, the open dots) shift away from those of the WX_2 layer ($\tau\mathbf{K}$, the solid dots). An electron and a hole can recombine if their momentum sum $\mathbf{k} + \mathbf{k}'$ equals the displacement vector between a $\tau\mathbf{K}$ and a $\tau\mathbf{K}'$ corner (the green arrows). (b) The interlayer Coulomb interaction (the wavy line) between the electron and the hole conserves their momentum sum [36]. An interlayer exciton formed with kinematic momentum $\mathbf{Q} = \pm\mathbf{Q}_0$, $\pm\hat{C}_3\mathbf{Q}_0$, or $\pm\hat{C}_3^2\mathbf{Q}_0$ can directly recombine to emit a photon. (c) These light cones at finite velocities are illustrated by the bright spots on the energy dispersion, while the exciton is optically dark at the energy minimum at $\mathbf{Q} = 0$.

the displacement from the MoX_2 Brillouin zone (BZ) corner $\tau\mathbf{K}'$ to a nearest WX_2 BZ corner $\tau\mathbf{K}$, with $\tau' = \tau$ ($\tau' = -\tau$) if the twist angle θ is near 0° (60°). Thus, in the phase space parametrized by \mathbf{Q} , there are six light cones at finite velocities [Fig. 1(c)]. With low symmetry of the exciton states, photon emission at each light cone can have a general form of being elliptically polarized. The exciton states at \mathbf{Q}_0 , $\hat{C}_3\mathbf{Q}_0$, and $\hat{C}_3^2\mathbf{Q}_0$ have the same electron-hole valley configuration (τ', τ) and are related by the \hat{C}_3 rotation, so their elliptical polarizations are also \hat{C}_3 rotations of each other. The rest of the three light cones are their time-reversal counterparts with the valley configuration $(-\tau', -\tau)$ [Fig. 2(a)].

With the helicity determined by the valley index and the major axis of the elliptical polarization locked to the exciton moving direction, the finite-velocity light cones make possible the injection of valley polarization as well as valley current [Fig. 2(a)]. Under excitation by a linearly polarized light, the absorption rates are all different at the

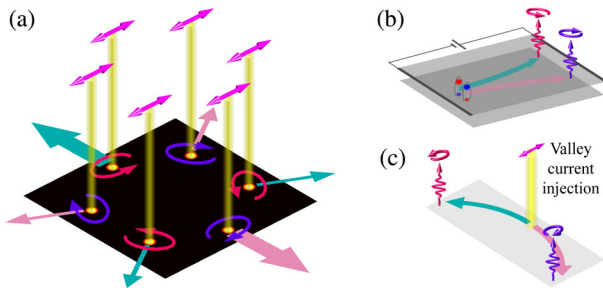


FIG. 2 (color online). (a) Injection of pure valley current of excitons by linearly polarized light (the green and pink colors denoting the valleys). See the text. (b) Valley Hall effect in heterobilayers. With the light cones at finite velocities, the spatial pattern of emission polarization also depends on the exciton distribution in \mathbf{Q} space. (c) The inverse valley Hall effect.

six light cones associated with different velocities, giving rise to valley current [Fig. 2(a)]. Using ϕ to denote the angle between the major axis of polarization of the light cone at $\mathbf{Q}_0 = Q_0\mathbf{y}$ and that of a linearly polarized pumping light, a pure valley current is injected with rate $\frac{3}{4}[1 - (1 - \epsilon)^2]/[1 + (1 - \epsilon)^2](\hbar K\theta/M_0)\alpha$ in the direction $\mathbf{x} \sin 2\phi + \mathbf{y} \cos 2\phi$, with α being the population injection rate.

This exciton system is an ideal platform to study the Berry phase effect in the Bloch band. The valley dependent Berry curvatures of electrons and holes in monolayer TMDs can be inherited by the exciton [6,44,45], giving rise to an exciton valley Hall effect that is observable from the spatial pattern of emission polarization [Fig. 2(b)] [44,46]. The elliptically polarized emission on the two edges with opposite helicity is a phenomenon unique to this heterobilayer interlayer exciton system. The optical injection of excitonic valley current through the finite-velocity light cones further makes possible the observation of the inverse valley Hall effect, where the valley current induces a number current in perpendicular direction due to the skewed motion by the Berry curvature [Fig. 2(c)].

We present now the theoretical details that back up the above general statements on the anomalous light cones. We first give the description of interlayer valley excitons at vanishing interlayer coupling, and we then analyze their light-interaction effects due to the residue interlayer coupling. Using $|e_{\tau', \mathbf{k}'}\rangle$ ($|h_{\tau, \mathbf{k}}\rangle$) to denote the electron (hole) Bloch state at momentum \mathbf{k}' (\mathbf{k}) measured from $\tau\mathbf{K}'$ ($\tau\mathbf{K}$) in MoX_2 (WX_2) layer in the *absence* of interlayer hybridization, we write the exciton wave function as [36,47]

$$|X_{\tau'\tau, \mathbf{Q}}^{(0)}\rangle = \sum_{\Delta\mathbf{Q}} \Phi_I(\Delta\mathbf{Q}) |e_{\tau', \frac{m_e\mathbf{Q} + \Delta\mathbf{Q}}{M_0}}\rangle |h_{\tau, \frac{m_h\mathbf{Q} - \Delta\mathbf{Q}}{M_0}}\rangle, \quad (1)$$

with the eigenenergy $E_I(\mathbf{Q}) = (\hbar^2 Q^2/2M_0) + \Delta_g - E_b$. $M_0 \equiv m_e + m_h$, Δ_g is the band gap and E_b the binding energy, and $\Phi_I(\Delta\mathbf{Q})$ describes the electron-hole relative motion with $\Delta\mathbf{Q} \equiv (m_h\mathbf{k}' - m_e\mathbf{k})/M_0$. $\mathbf{Q} \equiv \mathbf{k} + \mathbf{k}'$ is the kinematical momentum associated with the center-of-mass velocity $\langle X_{\mathbf{Q}}^{(0)} | [(m_e/M_0)\hat{\mathbf{r}}_e + (m_h/M_0)\hat{\mathbf{r}}_h] | X_{\mathbf{Q}}^{(0)} \rangle = (\hbar/M_0)\mathbf{Q}$ [36], a good quantum number even in the absence of translational invariance.

The electron (hole) Bloch function can be expanded by $e^{i(\mathbf{k}' + \tau\mathbf{K}')\cdot\mathbf{r}}$ ($e^{i(\mathbf{k} - \tau\mathbf{K})\cdot\mathbf{r}}$), where $\tau\mathbf{K}' \equiv \tau\mathbf{K}' + \mathbf{G}'$ ($\tau\mathbf{K} \equiv \tau\mathbf{K} + \mathbf{G}$), and \mathbf{G}' (\mathbf{G}) denotes all of the reciprocal lattice vectors in MoX_2 (WX_2). An electron-hole pair $|e_{\tau', \mathbf{k}'}\rangle |h_{\tau, \mathbf{k}}\rangle$ can recombine only when $\mathbf{k}' + \tau\mathbf{K}' = -\mathbf{k} + \tau\mathbf{K}$. So, the exciton light cones are at the \mathbf{Q} -space points:

$$\mathbf{Q} = \tau\mathbf{K} - \tau'\mathbf{K}'. \quad (2)$$

In Figs. 3(c) and 3(e), $\tau\mathbf{K}$ and $\tau'\mathbf{K}'$ are shown in the extended zone scheme for a $\text{MoSe}_2/\text{WSe}_2$ heterobilayer at $\theta = 5^\circ$ and a $\text{MoS}_2/\text{WSe}_2$ heterobilayer at $\theta = 0^\circ$. When θ changes, $\tau'\mathbf{K}'$ ($\tau\mathbf{K}$) move on concentric circles denoted

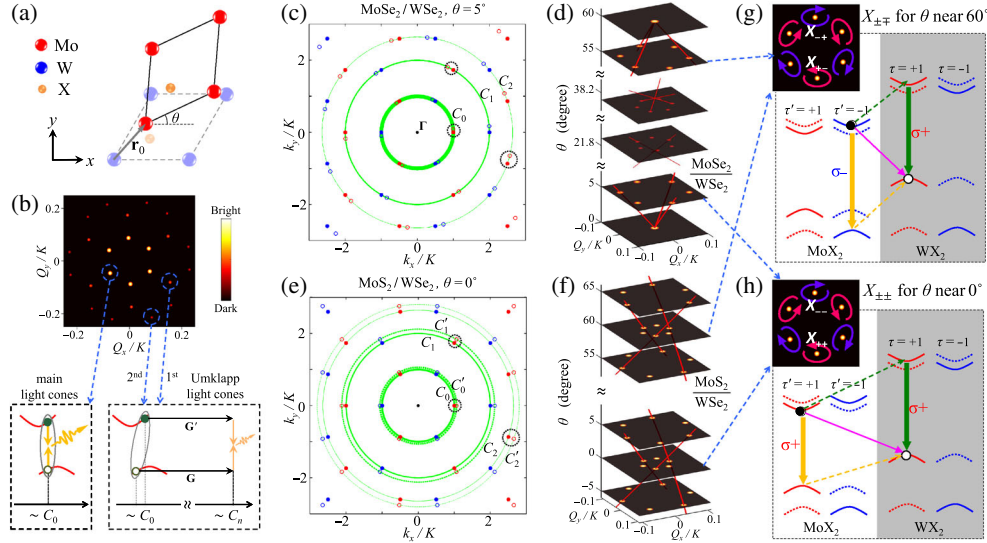


FIG. 3 (color online). (a) A MoX_2/WX_2 heterobilayer can be specified by the twist angle θ and the interlayer translation \mathbf{r}_0 . (b) The bright spots from large to small denote, respectively, the main, first-Umklapp, and second-Umklapp light cones in the \mathbf{Q} space, for $\text{WSe}_2/\text{MoSe}_2$ heterobilayers with θ near 0° or 60° . (c) Solid (open) dots denote $\tau\mathbf{k}$ ($\tau'\mathbf{k}'$) points of the WSe_2 (MoSe_2) layer in the extended zone scheme for $\theta = 5^\circ$. The locations of main (n th Umklapp) light cones are given by the displacements between the nearest solid and open dots on the C_0 (C_n) circles. (d) Low energy light cones at $\theta = 0^\circ, 5^\circ, 22.8^\circ, 37.2^\circ, 55^\circ, 60^\circ$, respectively. Evolutions of the main (second-Umklapp) light cones with θ are shown by the thick (thin) red lines. (e), (f) Lattice mismatched $\text{MoSe}_2/\text{WSe}_2$ heterobilayer. All light cones are always at a finite \mathbf{Q} . (g) Recombination pathways of interlayer exciton X_{-+} at θ near 60° . Solid (dashed) arrows denote the dipole transition (interlayer hopping). The inset shows the elliptical polarizations of emissions at the main light cones. (h) X_{++} at θ near 0° . At the finite \mathbf{Q} light cones, the polarization and strength of the optical dipole are independent of \mathbf{r}_0 .

as C'_n (C_n). The nearest $(\tau\mathbf{k}, \tau'\mathbf{k}')$ pairs on C_0 and C'_0 correspond to the main light cones. And the nearest $(\tau\mathbf{k}, \tau'\mathbf{k}')$ pairs on C_n and C'_n define the n th Umklapp light cones, where light coupling is an Umklapp-type process [Fig. 3(b)]. Figures 3(d) and 3(f) show the locations of these light cones as a function of twist angle θ , in a range $Q \leq 0.1K$ that corresponds to a kinetic energy $(\hbar^2 Q^2/2M_0) \lesssim 60$ meV. The six main light cones are located at $\mathbf{Q}_0 = (K - K' \cos \delta\theta)\mathbf{x} - K' \sin \delta\theta\mathbf{y}$ and its $n\pi/3$ rotations, where $\delta\theta$ is the deviation of θ from 0° or 60° . First-Umklapp light cones also have sixfold degeneracy, located at $2\mathbf{Q}_0$. The main and first-Umklapp light cones are the lowest energy ones at θ near 0° or 60° . There are 12 second-Umklapp light cones [Fig. 3(b)], and six of them become the lowest energy ones when θ is near the commensurate angle 21.8° or 38.2° in Fig. 3(d).

The interlayer hoppings (\hat{H}_T) will correct the wave function in Eq. (1), hybridizing $|X_{\tau'\tau, \mathbf{Q}}^{(0)}\rangle$ with small components of the intralayer valley excitons [36]. The interlayer exciton then acquires a finite transition dipole:

$$\begin{aligned} \mathcal{D}_{\tau'\tau, \mathbf{Q}} = & \langle 0 | \hat{\mathbf{D}} | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle + \langle 0 | \hat{\mathbf{D}} | X_{\tau'} \rangle \frac{\langle X_{\tau'} | \hat{H}_T | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle}{E_I(\mathbf{Q}) - E_M} \\ & + \langle 0 | \hat{\mathbf{D}} | X_{\tau} \rangle \frac{\langle X_{\tau} | \hat{H}_T | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle}{E_I(\mathbf{Q}) - E_W}. \end{aligned} \quad (3)$$

Here, $\hat{\mathbf{D}}$ is the electric dipole operator and $|0\rangle$ is the vacuum state. The first term on the right-hand side is from the transition dipole between $|e_{\tau', \mathbf{k}'}\rangle$ and $|h_{\tau, \mathbf{k}}\rangle$, which is small because of the spatial separation. The second and third terms describe, respectively, the light coupling mediated by bright intralayer exciton $X_{\tau'}$ in the MoX_2 layer through hole interlayer hopping, and X_{τ} in WX_2 through electron hopping. E_M and E_W are the energies of the intralayer excitons, a few hundred meV above E_I as determined from optical and scanning tunneling measurements [19–21, 30]. $\mathcal{D}_{\tau'\tau, \mathbf{Q}}$ can be extracted from *ab initio* calculations at several commensurate stackings [36], from which we can determine the dipole strength, the ellipticity, and the major axes of the emission polarization.

In the two-center approximation [48, 49],

$$\begin{aligned} \langle 0 | \hat{\mathbf{D}} | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle & \propto \sum_{\mathbf{k}, \mathbf{k}'} \delta_{\mathbf{Q}, \tau\mathbf{k} - \tau'\mathbf{k}'} \mathbf{D}_0^{2\tau}(\tau\mathbf{k}) e^{-i\tau\mathbf{k} \cdot \mathbf{r}_0}, \\ \langle X_{\tau'} | \hat{H}_T | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle & \approx \sum_{\mathbf{k}, \mathbf{k}'} \delta_{\mathbf{Q}, \tau\mathbf{k} - \tau'\mathbf{k}'} t_{2\tau'}^{2\tau}(\tau\mathbf{k}) e^{-i\tau\mathbf{k} \cdot \mathbf{r}_0}, \\ \langle X_{\tau} | \hat{H}_T | X_{\tau'\tau, \mathbf{Q}}^{(0)} \rangle & \approx \sum_{\mathbf{k}, \mathbf{k}'} \delta_{\mathbf{Q}, \tau\mathbf{k} - \tau'\mathbf{k}'} t_0^0(\tau\mathbf{k}) e^{-i\tau\mathbf{k} \cdot \mathbf{r}_0}, \end{aligned} \quad (4)$$

where $t_{m'}^m$ ($\mathbf{D}_{m'}^m$) are the Fourier components of the hopping integral (transition dipole) between W and Mo d orbitals with magnetic quantum numbers m and m' , respectively, which decay fast with the increase of $|\mathbf{k}|$ [36]. Equation (4)

shows that the strength and polarization of $\mathcal{D}_{\tau'\tau}$ at the sixfold light cones have no dependence on the layer translation \mathbf{r}_0 , which appears in a phase factor $e^{-i\tau\mathbf{k}\cdot\mathbf{r}_0}$ only.

Figure 3(h) schematically shows interlayer excitons $X_{\tau'\tau}$ with $\tau' = \tau$. These are the ones in the main and first-Umklapp light cones at θ near 0° and in second-Umklapp light cones at θ near 38.2° . Figure 3(g) shows $X_{\tau'\tau}$ with $\tau' = -\tau$, which are the ones in the main and first-Umklapp light cones at θ near 60° and in second-Umklapp light cones at θ near 21.8° . The emission polarizations of the main light cones are illustrated in the insets. For MoSe₂/WSe₂ heterobilayers, we find the emission polarization of X_{++} at θ near 0° has an ellipticity $\epsilon = 0.35$, which correspond to a ratio of 5 between the $\mathbf{e}_+ \equiv (\mathbf{x} + i\mathbf{y})/\sqrt{2}$ and $\mathbf{e}_- \equiv (\mathbf{x} - i\mathbf{y})/\sqrt{2}$ components of $\mathcal{D}_{++,\mathbf{Q}}$. The \mathbf{e}_+ component is contributed by all three terms in Eq. (3), while the much smaller \mathbf{e}_- component is from the $\langle 0|\hat{\mathbf{D}}|X_{\tau'\tau,\mathbf{Q}}^{(0)}\rangle$ term only. X_{--} is simply the time reversal of X_{++} . At θ near 60° , the valley configurations X_{+-} and X_{-+} have $\epsilon = 0.33$.

From *ab initio* calculations [36], $X_{\pm\pm}$ at θ near 0° has a transition dipole strength of $\sim 0.052D$ at the main light cones, and $\sim 0.001D$ at both the first- and second-Umklapp light cones. $X_{\pm\mp}$ at θ near 60° has a transition dipole strength of $\sim 0.052D$ at the main light cones, $\sim 0.001D$ at first-Umklapp, and $\sim 0.0005D$ at second-Umklapp light cones. D is the dipole strength of the intralayer excitons [36,50].

For MoSe₂/WSe₂ or MoS₂/WS₂ heterobilayers, neglecting the $\sim 0.1\%$ lattice mismatch, they have commensurate stacking at $\theta = 0^\circ, 60^\circ, 21.8^\circ$, or 38.2° . At $\theta = 0^\circ, 60^\circ$, all main and Umklapp light cones merge into a single cone at $\mathbf{Q} = 0$. At $\theta = 21.8^\circ, 38.2^\circ$, the six second-Umklapp light cones merge into a single one at $\mathbf{Q} = 0$ [cf. Fig. 3(d)]. For the light cone at $\mathbf{Q} = 0$, the transition dipole is the superposition of multiple $(\tau\mathbf{k}, \tau'\mathbf{k}')$ pairs, each of which is associated with a distinct phase factor $e^{-i\tau\mathbf{k}\cdot\mathbf{r}_0}$ [cf. Eq. (4)]. Their interference then gives rise to the different \mathbf{r}_0 dependence for the $\sigma+$ and $\sigma-$ components of $\mathcal{D}_{\tau'\tau}$ [36]:

$$\begin{aligned} \mathbf{e}_\tau \cdot \mathcal{D}_{\tau'\tau} &\propto e^{-i\tau\mathbf{K}\cdot\mathbf{r}_0} + e^{-i\tau\hat{C}_3\mathbf{K}\cdot\mathbf{r}_0} + e^{-i\tau\hat{C}_3^2\mathbf{K}\cdot\mathbf{r}_0}, \\ \mathbf{e}_{-\tau} \cdot \mathcal{D}_{\tau'\tau} &\propto e^{-i\tau\mathbf{K}\cdot\mathbf{r}_0} + e^{-i\tau[\hat{C}_3\mathbf{K}\cdot\mathbf{r}_0 + (2\pi/3)]} + e^{-i\tau[\hat{C}_3^2\mathbf{K}\cdot\mathbf{r}_0 + (4\pi/3)]}. \end{aligned} \quad (5)$$

Thus, both the strength and the elliptical polarization of the transition dipole change with \mathbf{r}_0 [51], as shown in Fig. 4. The above relations dictate that at $\mathbf{r}_0 = 0$ and $(\mathbf{a}_1 + \mathbf{a}_2)/3$, the in-plane transition dipole of $X_{\tau'\tau}$ becomes fully circularly polarized, while at $\mathbf{r}_0 = 2(\mathbf{a}_1 + \mathbf{a}_2)/3$ both \mathbf{e}_\pm components of $\mathcal{D}_{\tau'\tau}$ vanish, in agreement with the *ab initio* calculations in Fig. 4 [54]. Other values of \mathbf{r}_0 break the threefold rotational symmetry, so elliptically polarized emission by $X_{\tau'\tau}$ at $\mathbf{Q} = 0$ is allowed. For $\theta = 0^\circ$, the transition dipole is strongest $\sim 0.16D$ at $\mathbf{r}_0 = 0$

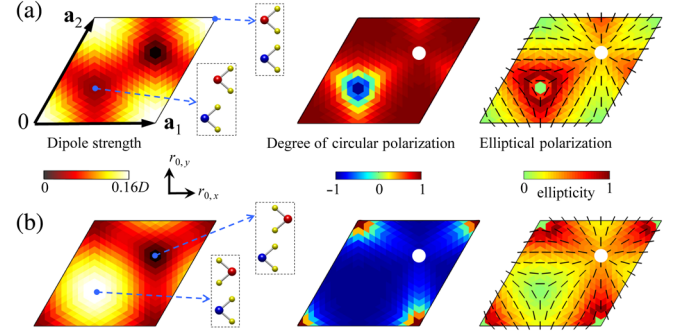


FIG. 4 (color online). Dependence of the transition dipole strength and polarization as a function of layer translation \mathbf{r}_0 in lattice matching MoSe₂/WSe₂ heterobilayers. (a) X_{++} at AA-like stacking ($\theta = 0^\circ$). The three panels show, respectively, the dipole strength, the degree of polarization in circular basis, and the ellipticity and major axis (represented by short lines) of elliptical polarization, from *ab initio* calculations [54]. (b) The same plots for X_{-+} at AB-like stacking ($\theta = 60^\circ$).

[cf. Fig. 4(a)], while for $\theta = 60^\circ$, the transition dipole is strongest $\sim 0.15D$ at $\mathbf{r}_0 = (\mathbf{a}_1 + \mathbf{a}_2)/3$ [cf. Fig. 4(b)].

For interlayer excitons relaxed to the bottom of the energy dispersion, their photon emission need the assistance of phonon or impurity scattering into the finite-velocity light cones. The inefficiency of such a process, compared to the direct recombination of intralayer excitons, can be relevant to the long lifetime observed in the time resolved photoluminescence (PL) [19], which can be interesting for future study. The energy of the emitted photon has contributions from the exciton kinetic energy and the energy of the phonon (if involved), the distribution of which may account for the large broadening of PL. Our results also explain why strong interlayer exciton PL is observed only at θ near 0° or 60° [19], where the main light cones get sufficiently close to the bottom of energy dispersion. At θ near 21.8° or 38.2° the second-Umklapp light cones are close to $\mathbf{Q} = 0$; however, the weak dipole of these light cones renders the PL inefficient.

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*wangyao@hku.hk

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