

# Kinetically Controlled Self-Assembly of Phosphorescent Au<sup>III</sup> Aggregates and Ligand-to-Metal-Metal Charge Transfer Excited State: A Combined Spectroscopic and DFT/TDDFT Study

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**ABSTRACT:** Metallophilic interactions in d<sup>10</sup>-d<sup>10</sup>(Au<sup>I</sup>-Au<sup>I</sup>)/d<sup>8</sup>-d<sup>8</sup>(Pt<sup>II</sup>-Pt<sup>II</sup>, Rh<sup>I</sup>-Rh<sup>I</sup>, Ir<sup>I</sup>-Ir<sup>I</sup>) complexes have been widely studied for decades, and metal-metal (M-M) bonding character has been revealed in both the ground and excited states. These M-M closed-shell interactions are appealing driving forces for the self-assembly of supramolecular/polymeric systems, providing luminescent properties distinctly different from those of the corresponding monomer. However, reports on attractive interactions between two Au<sup>III</sup> complex cations are scarce in the literature. Herein is described a series of pincer-type cationic Au<sup>III</sup> complexes with different auxiliary ligands, among which the Au<sup>III</sup>-allenylidene complex displays a close Au-Au contact of 3.367 Å between neighboring molecules in its X-ray crystal structure; Au<sup>III</sup>-isocyanide complexes show a broad redshifted absorption band and prominent phosphorescence upon aggregation that was influenced by an attractive Au<sup>III</sup>-Au<sup>III</sup> bonding interaction in the excited state; and Au<sup>III</sup>-acetylene complexes can undergo living supramolecular polymerization upon varying the counteranion. The nature of the emissive excited state(s) of the Au<sup>III</sup> aggregates is assigned to a mixture of major <sup>3</sup>[π - π\*] and minor <sup>3</sup>LMMCT (ligand-to-metal-metal charge transfer) states based on combined spectroscopic and DFT/TDDFT studies. The morphology of the Au<sup>III</sup> aggregates is highly dependent on the concentration and nature of the counteranion. A qualitative model has been applied to account for the concentration- and counteranion-dependent kinetics of the supramolecular polymerization process.

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## Introduction

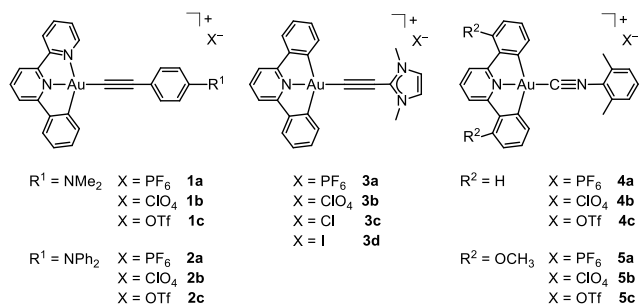
Aurophilic attraction, commonly referred to Au<sup>I</sup>-Au<sup>I</sup> interaction, has been subjected to extensive experimental and theoretical studies for decades.<sup>1,2</sup> Pioneering works on Au<sup>I</sup> systems<sup>2,3</sup> by Pyykkö and coworkers suggest that aurophilic attraction arises mainly from dispersion and virtual charge transfer interactions. In contrast to aurophilicity in Au<sup>I</sup> chemistry and d<sup>10</sup>-d<sup>10</sup> interactions,<sup>2,4-7</sup> investigations on d<sup>8</sup>-d<sup>8</sup> Au<sup>III</sup>-Au<sup>III</sup> interactions are sparse.<sup>8,9</sup> Whether Au<sup>III</sup> complexes would display specific low energy electronic transitions upon aggregation, similar to their isoelectronic Pt<sup>II</sup> and Pd<sup>II</sup> counterparts, also remains an unanswered question. For decades, square planar Pt<sup>II</sup> complexes have been widely studied with regard to aspects related to <sup>1/3</sup>MMLCT excited states, supramolecular self-assembly, and photoluminescence, all of which are impacted by the presence of closed-shell Pt<sup>II</sup>-Pt<sup>II</sup> interactions.<sup>10-15</sup> Recent studies showed that Pd<sup>II</sup> complexes can also display Pd<sup>II</sup>-Pd<sup>II</sup> interactions and <sup>1/3</sup>MMLCT excited states upon aggregation.<sup>16,17</sup> A series of photoresponsive semiconductors and sensors based on pincer-type Pt<sup>II</sup> complexes could also be fabricated.<sup>10,18-20</sup> As Au<sup>III</sup> ion has a higher redox potential and is more electrophilic than Pt<sup>II</sup> ion, supramolecular materials with new functions could be obtained through a self-assembly process using Au<sup>III</sup> complexes as building blocks.

Recently, kinetically controlled supramolecular polymerization of cationic Pt<sup>II</sup> and Pd<sup>II</sup> complexes has been achieved by modulating the pathway complexity.<sup>21,22</sup> Besenius and coworkers reported the stepwise self-assembly of Au<sup>I</sup>-metallopeptides in water by making use of Au<sup>I</sup>-Au<sup>I</sup> interac-

tions.<sup>23</sup> However, to the best of our knowledge, pincer Au<sup>III</sup> complexes that undergo a kinetically controlled supramolecular polymerization process are unprecedented, and the excited states of Au<sup>III</sup> aggregates are scarcely studied.

Kinetic control over supramolecular polymerization, which is also very ubiquitous in the synthesis of biological structures, has been found to be a powerful tool in preparing highly ordered nanostructures with tailored and specific functions.<sup>24-31</sup> An appropriate kinetic energy barrier is the key factor for molecules undergoing a supramolecular polymerization process out of their thermodynamic equilibrium.<sup>32</sup> Seeded supramolecular polymerization can be achieved by manipulating the pathway complexity; however, until recently, only limited studies in this area have been reported.<sup>21,24,26,27,32-42</sup>

Herein, compelling evidence that the pincer-type Au<sup>III</sup> complexes undergo kinetically controlled self-assembly, which can be devised to follow the living supramolecular polymerization, is described. A qualitative kinetic model<sup>43</sup> was applied to explain the pathway complexity of the supramolecular self-assembly of pincer Au<sup>III</sup> complexes. Excited states with mixed <sup>1/3</sup>LMMCT and <sup>1/3</sup>[π - π\*] character have been assigned to account for the spectroscopic properties of Au<sup>III</sup> aggregates based on DFT/TDDFT calculations and experimental observations.



**Scheme 1. Chemical Structures of Au<sup>III</sup> Complexes.**

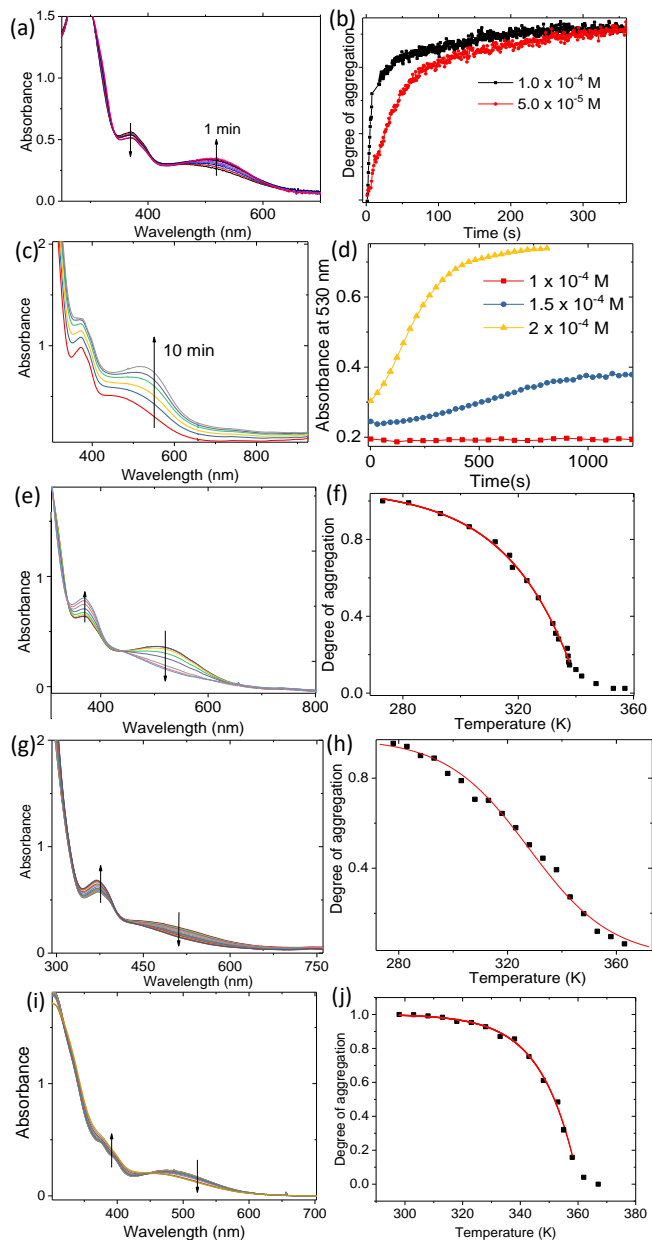
## Results

Complexes **1a-b**, **2a-c** were synthesized by the reactions between [Au(C<sup>^N^N</sup>)OAc]<sup>+</sup> and the corresponding acetylene according to the reported procedures with modifications;<sup>8,9</sup> among these compounds, complex **1a** has been reported with its X-ray crystal structure. In the crystal structure, complex **1a** was observed to be packed in a head-to-tail style, and the Au-Au distance is characterized to be approximately 3.49 Å. Complexes **3a-d** were synthesized according to the literature<sup>17</sup> procedure for preparing a series of pincer-type Au<sup>III</sup> complexes with different allenyldiene ligands. Complexes **4a-c**, **5a-c** were synthesized by treatment of [Au(C<sup>^N^C</sup>)Cl] with 2,6-dimethylphenyl isocyanide and an excess of salt (NH<sub>4</sub>PF<sub>6</sub> or KPF<sub>6</sub>, LiClO<sub>4</sub>, AgOTf or LiOTf) in dry CH<sub>3</sub>CN solution under Ar for 12 h. Complexes **4a-c** are relatively stable in the solid state (gradually decomposed after approximately 1 week under air; **4c** is the most unstable one), but they are highly sensitive to water in the solution state and quickly decomposed when water is added to their acetonitrile solution (within ten minutes). Compared to **4a-c**, complexes **5a-c** with electron-donating OCH<sub>3</sub> substituents on the [C<sup>^N^C</sup>] ligand scaffold show better stability in both solution and the solid state. In a water and acetonitrile solution mixture, both **5a** and **5c** are stable for several hours; complex **5b** is stable in solution state with no decomposition within one week. For the unreported stable Au<sup>III</sup> complexes (**1b-2c**, **3a-d**) studied in this work, NMR spectroscopy, high-resolution mass spectrometry and element analysis were used for characterization. For relatively unstable complexes **4a-5c**, NMR spectroscopy and high-resolution mass spectrometry were combined to characterize these complexes. Furthermore, the X-ray crystal structure of complex **3d** was determined.

## Supramolecular Polymerization of Complexes 1 and 2

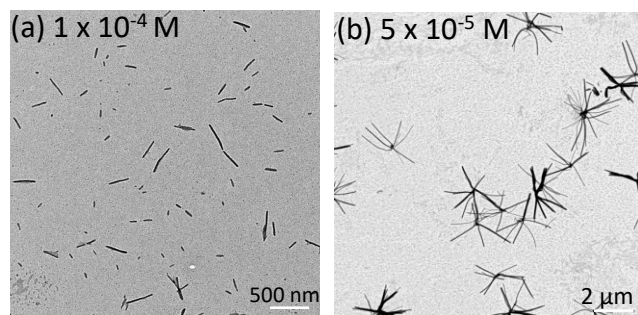
Complexes **1a-c**, **2a-c** in CH<sub>3</sub>CN (1 × 10<sup>-4</sup> M) display moderate absorption bands at approximately 400-600 nm (Figure S1). When a mixed solvent of H<sub>2</sub>O/CH<sub>3</sub>CN was used, a broad redshifted absorption band appeared, indicating the formation of aggregate species. A kinetically controlled supramolecular polymerization process was revealed for **1a** and **1b** in the mixed solvent of H<sub>2</sub>O and CH<sub>3</sub>CN (9:1 v/v). In Figure 1a-d, a new broad band at approximately 530 nm gradually appeared for **1a** and **1b** in the time course of the experiment. Decreasing the complex concentration slowed down this process. The lag time shown in Figure 1d (1.5 × 10<sup>-4</sup> M) for **1b** indicates the existence of a metastable aggregate state that prohibited spontaneous aggregation. By changing the counteranion from ClO<sub>4</sub><sup>-</sup> (**1b**) to PF<sub>6</sub><sup>-</sup> (**1a**), the lag time, indicating the presence of a kinetically metastable state, disappeared (black line in Figure 1b at a concentration of 1 × 10<sup>-4</sup> M). We propose that the energy barrier between the

metastable and thermodynamic aggregates is much higher for **1b** than for **1a**. As a result, the high kinetic barrier of **1b** seriously retarded the transformation between these two aggregate states.

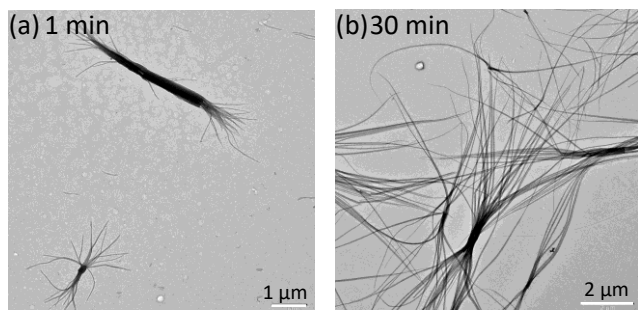


**Figure 1.** (a) Time-dependent UV/Vis absorption spectral changes of **1a** (1 × 10<sup>-4</sup> M, H<sub>2</sub>O/CH<sub>3</sub>CN, 9:1, v/v) at 298 K. (b) Time-dependent degree of aggregation of **1a** calculated from the absorbance at 530 nm at the indicated concentration. (c) Time-dependent UV/Vis absorption spectral changes of **1b** (2 × 10<sup>-4</sup> M, H<sub>2</sub>O/CH<sub>3</sub>CN, 9:1, v/v). (d) Time-dependent absorbance change of **1b** at 530 nm at the indicated concentration. (e, g, i) Temperature-dependent UV/Vis absorption spectral changes of (e) **1a**, (g) **1c** and (i) **2a** upon heating at a concentration of 1 × 10<sup>-4</sup> M (H<sub>2</sub>O/CH<sub>3</sub>CN, 9:1, v/v). (f, h, j) Plot of the degree of aggregation (black dot) as a function of temperature and the fitting curve (red line) of (f) **1a**, (h) **1c** and (j) **2a**.

Under the same conditions, no time-dependent UV/Vis absorption changes appeared for complexes **1c** and **2a-c**. Upon increasing the temperature, changes in the absorbance for **1a** at 530 nm (thermodynamic aggregate) are indicative of a cooperative mechanism (Figure 1e-f) for the self-assembly process. By fitting the absorption spectral changes with a cooperative model,<sup>29</sup> an elongation temperature of  $T_e = 341$  K and an enthalpy release of  $\Delta H = -9.9$  kcal/mol were obtained. For **1c**, the changes in the UV/Vis absorbance at 500 nm upon changing the temperature could be explained by the isodesmic supramolecular polymerization mechanism. A melting temperature of  $T_m = 328$  K and an enthalpy release of  $\Delta H = -15.4$  kcal/mol were obtained by fitting the absorption changes (Figure 1g-h) with the isodesmic model. Upon changing the substituent group from  $\text{NMe}_2$  (**1a**) to  $\text{NPh}_2$  (**2a**) in the acetylene ligand, a greater entropy release of  $\Delta H = -20.9$  kcal/mol was obtained, by fitting the temperature-dependent degree of aggregation to the cooperative self-assembly model (Figure 1i-j). The relatively large non-covalent bond formation energy indicates less dynamic features of **2a** (compared with that of **1a**) and therefore, disfavoring **2a** undergoing kinetically controlled self-assembly process.



**Figure 2.** TEM images of **1a** at (a)  $1 \times 10^{-4}$  M and (b)  $5 \times 10^{-5}$  M in  $\text{H}_2\text{O}/\text{CH}_3\text{CN}$  (9:1 v/v).



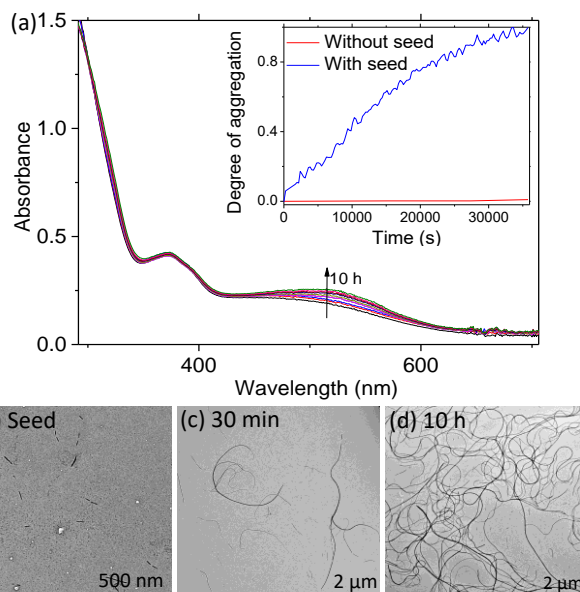
**Figure 3.** Time-dependent TEM changes of **1b** at (a) 1 min and (b) 30 min ( $1 \times 10^{-4}$  M,  $\text{H}_2\text{O}/\text{CH}_3\text{CN}$ , 9:1 v/v).

As the concentration of a complex can make a significant difference in the kinetics of its supramolecular polymerization, we assessed the nano-morphologies of **1a** by TEM images at different concentrations (Figures 2 and 3). When the self-assembly process of **1a** was conducted at a relatively high concentration ( $1 \times 10^{-4}$  M, Figure 2a), short 1D nanowire structures were observed. In contrast, when the experiment was performed at a relatively low concentration ( $5 \times 10^{-5}$  M, Figure 2b), two-dimensional (2D) nano-star-shaped structures were observed. The structures featured a nucleus

with diameters of approximately 100-200 nm and elongated polymers with a length of approximately  $2 \mu\text{m}$ . The morphologies of the supramolecular polymers were found to be highly dependent on the counteranion. We monitored the time-dependent morphology change in **1b** at a concentration of  $1 \times 10^{-4}$  M (Figure 3). In the TEM images, short nanowires emerged in the initial stage. Gradually, thin and long nanowires began to grow at the terminals of the nucleus, leading to 2D bowknot-shaped structures.

### Living Supramolecular Polymerization of Complex **1**

The supramolecular polymerization of **1b** at a low concentration (below  $1 \times 10^{-4}$  M) was monitored by UV/Vis absorption spectroscopy, and the transformation from kinetic aggregates to thermodynamic aggregates was found to be very slow (red line in Figure 1d at a concentration of  $1 \times 10^{-4}$  M). The long lag time is tentatively attributed to the high energy barrier between the thermodynamic self-assemblies and kinetic self-assemblies; however, such a barrier can be bypassed by adding a seed. Therefore, we prepared a seed from **1a** ( $2 \times 10^{-5}$  M) to facilitate the direct aggregation of **1b** ( $5 \times 10^{-5}$  M) into a final stable aggregate state during the supramolecular polymerization process.



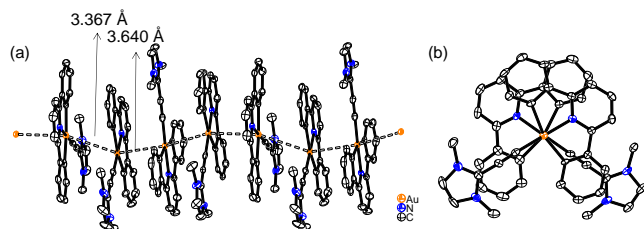
**Figure 4.** (a) Time-dependent UV/Vis absorption spectra of **1b** with the aid of a seed prepared from **1a** (**1b**,  $5 \times 10^{-5}$  M; seed **1a**,  $2 \times 10^{-5}$  M). Inset: degree of aggregation calculated at 530 nm for **1b** with and without the seed in the time course of the experiment. (b) TEM image of the seed. (c) and (d) Time-dependent TEM images of **1b** with the seed obtained at 30 min and 10 h, respectively.

As depicted in Figure 4a, after adding the seed to an acetonitrile solution of **1b** at  $5 \times 10^{-5}$  M, the lag time disappeared, indicating an externally seeded self-assembly growth process. The length of the seed was approximately 200 nm, as shown in Figure 4b. The time-dependent evolution of the aggregation process was recorded by following the TEM images, as depicted in Figure 4c-d. According to our previous work on counteranion-modulated supramolecular polymerization processes and computational studies,<sup>16,22,44</sup> the  $\text{PF}_6^-$  counteranion displays weaker intermolecular interactions than the  $\text{ClO}_4^-$  counteranion with cationic complexes. In

this work, upon breaking the Au<sup>III</sup> complex cation-PF<sub>6</sub><sup>-</sup> interaction, aggregates with different packing forms are proposed to be formed, which was mainly driven by the “cation-cation” dispersive interaction. For the seeded supramolecular polymerization process of **1b**, the use of **1a** (PF<sub>6</sub><sup>-</sup> counteranion) as a seed helped break the Au<sup>III</sup> complex cation-ClO<sub>4</sub><sup>-</sup> interaction, thereby facilitating the growth of its thermodynamic aggregates. To further prove the external seeded supramolecular polymerization, we performed two control experiments, monitoring the self-assembly of complex **1b** without seed **1a** and seed **1a** in the absence of **1b** by UV/Vis absorption spectroscopy (Figure S3).

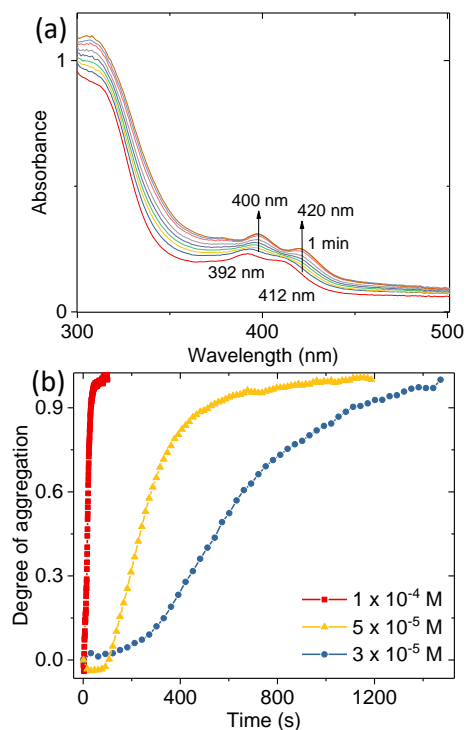
### Supramolecular Polymerization of Complex **3** and X-ray Crystal Structure

Organometallic complexes containing N-heterocyclic allenylidene (C=C=CRR') ligands are usually stable and exhibit intriguing photoluminescence, metallophilic interactions and cytotoxicity properties.<sup>45</sup> Complexes **3a-d** were prepared with N-heterocyclic allenylidene ligands, and these complexes are stable in both the solid state and solution. The crystal structure of **3d** shows a zig-zag Au-Au chain structure, featuring a close Au-Au contact of approximately 3.367 Å, with complex cations arranged in a skewed fashion, as shown in Figure 5.

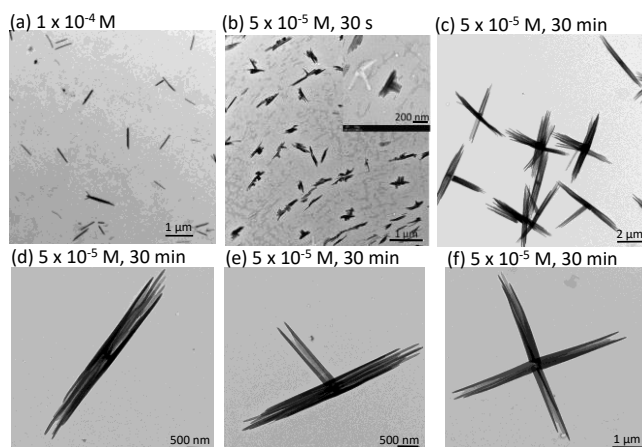


**Figure 5.** (a) X-ray crystal structure of **3d** featuring a zig-zag structure with close alternating Au-Au contacts of 3.367 Å and 3.640 Å. Hydrogen atoms and counteranions are omitted for clarity. (b) Top view of the dimeric Au<sup>III</sup>-Au<sup>III</sup> structure.

Complexes **3a-d** in CH<sub>3</sub>CN (1 × 10<sup>-4</sup> M), which are presumably in a discrete monomeric form, display two moderate absorption bands at approximately 380 and 400 nm (Figure S5-S6). When a mixed solvent of H<sub>2</sub>O/CH<sub>3</sub>CN was used, two redshifted absorption bands at approximately 400 nm and 420 nm appeared for **3a**, indicating the formation of aggregate species (Figure 6a). Time-dependent UV/Vis absorption spectral changes and stepwise self-assembly process were observed for this complex. As shown in Figure 6a, upon injecting water into an CH<sub>3</sub>CN solution of **3a**, two new absorption bands at approximately 400 and 420 nm gradually emerged; this behavior is indicative of the presence of a kinetically trapped aggregate state during the self-assembly pathway. Lowering the concentration of **3a** slowed down this process. As shown in Figure 6b, lag times of ~100 s and ~300 s were observed when the measurements were conducted at concentrations of 5 × 10<sup>-5</sup> M and 3 × 10<sup>-5</sup> M, respectively. For aggregates **3b-d** in the mixed solvent of H<sub>2</sub>O/CH<sub>3</sub>CN, absorption bands at around 392 nm and 412 nm appeared instantly with no time-dependent changes in Figure S5-S6. The absorption spectra of aggregates **3b-d** are similar to that observed for **3a**-kinetic aggregate in Figure 6a.



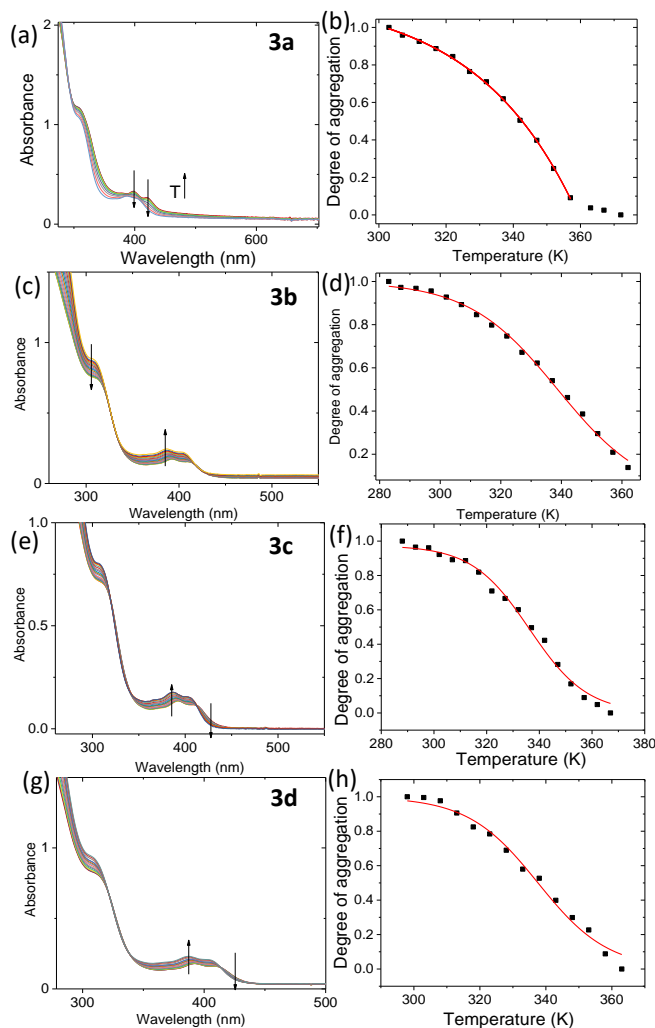
**Figure 6.** (a) Time-dependent UV/Vis absorption spectral changes of **3a** (1 × 10<sup>-4</sup> M, H<sub>2</sub>O/CH<sub>3</sub>CN = 9:1, v/v) at 298 K. Inset: amplification of the change in the absorbance at 400 and 420 nm in the time course of the experiment. (b) Degree of aggregation of **3a** calculated from the absorbance at 420 nm at the indicated concentration in the time course of the experiment.



**Figure 7.** TEM images of **3a** at (a) 1 × 10<sup>-4</sup> M in H<sub>2</sub>O/CH<sub>3</sub>CN (9:1 v/v). TEM images of **3a** at 5 × 10<sup>-5</sup> M in H<sub>2</sub>O/CH<sub>3</sub>CN (9:1 v/v) obtained at (b) 30 s and (c-f) 30 min.

The morphology of the nanostructures is highly dependent on the concentration of **3a** in the mixed CH<sub>3</sub>CN/H<sub>2</sub>O solution (Figure 7). When the self-assembly process was conducted at a concentration of 1 × 10<sup>-4</sup> M, a 1D nanowire morphology (with a length of approximately 1 μm) was observed. However, upon lowering the concentration to 5 × 10<sup>-5</sup> M, polymers characterized by a nucleus and elongated wires emerged in a stepwise manner. Changes in time-dependent

TEM images were monitored and captured: initially, a small nucleus with a length of  $\sim 200$  nm appeared (Figure 7b), followed by the Au<sup>III</sup> complexes growing at the terminals of the nucleus, leading to 2D supramolecular polymers (with a size of approximately 3-4  $\mu$ m), as shown in Figure 7d-f. By replacing the PF<sub>6</sub> counteranion with other counteranions, small nanoparticle structures (with a diameter of approximately 50 nm) were observed for **3b-d** and showed no obvious relationship with the complex concentration upon aggregation (Figures S7-S8). In its aggregate form, complex **3d** showed weak phosphorescence with an emission maximum at  $\sim 500$  nm, as shown in Figure S9.

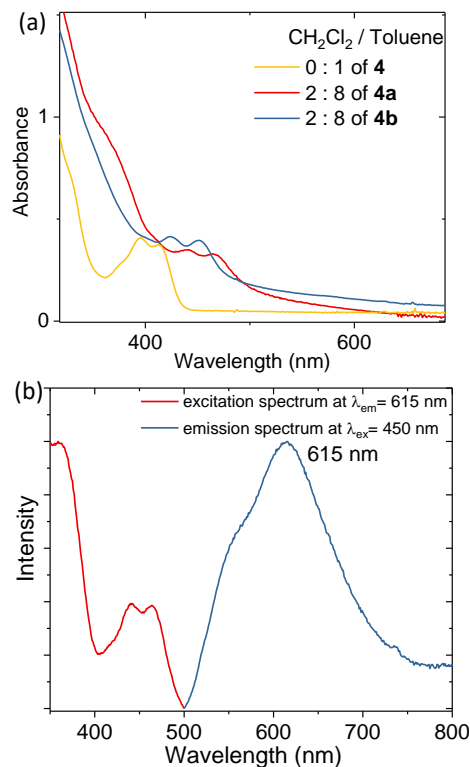


**Figure 8.** Temperature-dependent UV/Vis absorption spectral changes of (a) **3a**, (c) **3b**, (e) **3c** and (g) **3d** at the concentration of  $1 \times 10^{-4}$  M in H<sub>2</sub>O/CH<sub>3</sub>CN (9/1, v/v) upon heating. Plot of the degree of aggregation (black dot) as a function of temperature and the fitting curve (red line) of (b) **3a**, (d) **3b**, (f) **3c** and (h) **3d**.

We performed temperature-dependent UV/Vis measurements for **3a-d** in the mixed H<sub>2</sub>O/CH<sub>3</sub>CN solvent, to probe the difference of their supramolecular polymerization mechanism by changing counteranions. As shown in Figures 8a-b, the changes in the degree of aggregation for **3a** could be fitted with the cooperative nucleation-elongation model, from which an elongation temperature of  $T_e = 359$  K and an enthalpy release of  $\Delta H = -8.5$  kcal/mol were obtained. In con-

trast, complex **3b-d** followed sigmoidal isodesmic self-assembly model as shown in Figures 8c-h. The non-covalent bond formation entropy and melting temperature were determined to be  $-17.2$  kcal/mol and 335 K for **3b**,  $-21.8$  kcal/mol and 335 K for **3c**,  $-22.5$  kcal/mol and 338 K for **3d**, respectively. The relatively large non-covalent bond formation energy for **3b-d** (compared with that of **3a**) during the aggregation process suggests their less dynamic features, which is considered to be one of the reasons that leads to the absence of a kinetically controlled self-assembly process for **3b-d**.

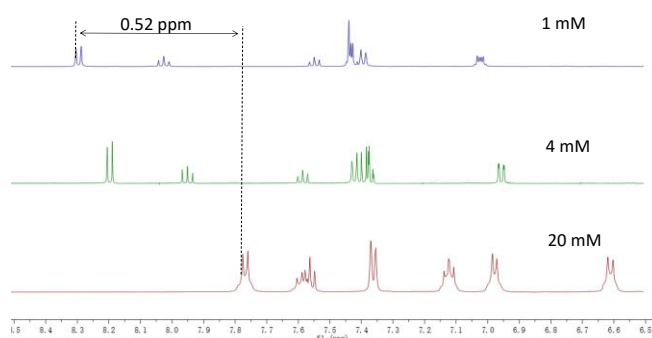
### Phosphorescent Au<sup>III</sup> Aggregates of Complexes 4 and 5



**Figure 9.** (a) Solvent-dependent UV/Vis absorption spectra of **4a** and **4b** ( $1 \times 10^{-4}$  M) at 298 K. The absorption spectra on **4a**, **4b** and **4c** in a monomeric state in CH<sub>2</sub>Cl<sub>2</sub> showed no differences. (b) Emission and excitation spectra of **4a** ( $1 \times 10^{-4}$  M, CH<sub>2</sub>Cl<sub>2</sub>/toluene = 2/8, v/v) at 298 K.

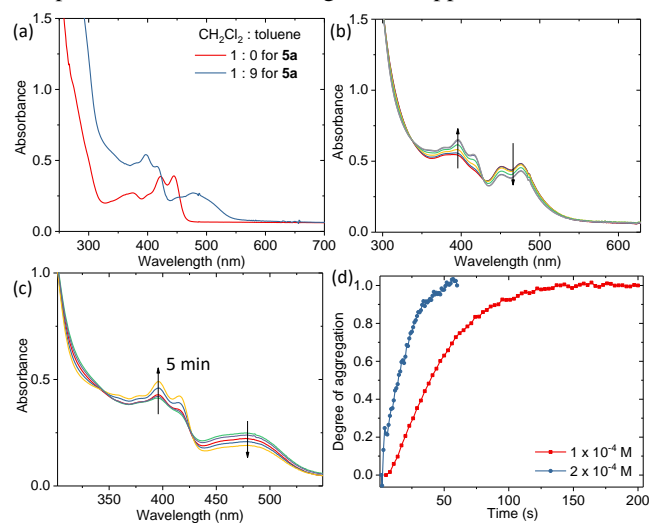
Au<sup>III</sup> complexes **4a-c** with isocyanide auxiliary ligands are sensitive to water and were observed to decompose after long-term standing under air in the solution state (Figures S10-S11). Therefore, we used dry toluene to induce supramolecular polymerization in CH<sub>2</sub>Cl<sub>2</sub> solution. The UV/Vis absorption spectrum of **4** in CH<sub>2</sub>Cl<sub>2</sub> (presumably in a discrete monomeric state) is shown in Figure 9a, together with the spectrum recorded in mixed CH<sub>2</sub>Cl<sub>2</sub>/toluene (2:8, v/v) corresponding to the aggregate state. Complex **4** in CH<sub>2</sub>Cl<sub>2</sub> solution ( $1 \times 10^{-4}$  M) shows moderate absorption bands at 395 and 415 nm. Upon aggregation, two new absorption bands of the aggregate species appeared at 420 and 445 nm for **4b** and at 443 and 470 nm for **4a**. Complex **4a** was found to be emissive in the aggregate state ( $1 \times 10^{-4}$  M, CH<sub>2</sub>Cl<sub>2</sub>/toluene, 2:8, v/v), featuring a broad emission peak at 615 nm, as shown in Figure 9b. For **4c**, a new absorption band at approximately

450 nm gradually appeared, and the polymerization rate was found to be related to the complex concentration, as shown in Figure S13.



**Figure 10.** Concentration-dependent  $^1\text{H}$  NMR spectra for **5a** in  $\text{CD}_3\text{CN}$  at 298 K.

Compared to **4a-c**,  $\text{Au}^{\text{III}}$  complexes **5a-c** show better stability in both the solid and solution state. In concentration- (Figure 10) and temperature-dependent  $^1\text{H}$  NMR spectra (Figure S17) of **5a**, an upfield change in the chemical shifts was observed upon increasing the concentration or lowering the temperature. The formation of  $\pi$ - $\pi$  stacked aggregate species accounted for the change in chemical shifts, which were influenced by the ring current of the aromatic groups.<sup>46,47</sup> In Figure 10, it can be noted that upon increasing the concentration from 1 mM to 20 mM, complex **5a** showed an upfield chemical shift change of 0.52 ppm.

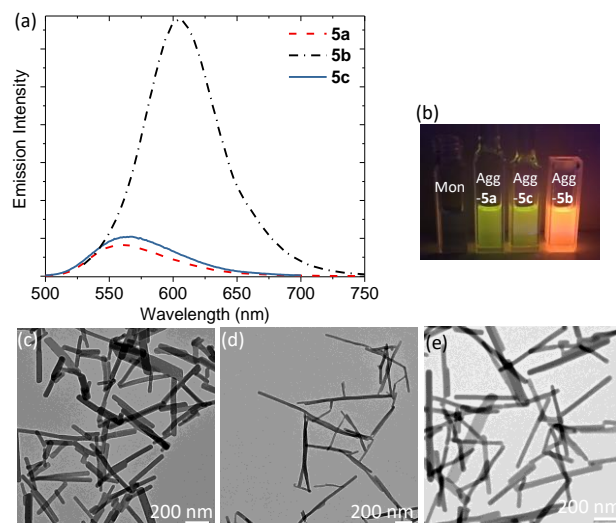


**Figure 11.** (a) Solvent-dependent UV/Vis absorption spectral changes for **5a** ( $1 \times 10^{-4}$  M) at 298 K. (b) Time-dependent UV/Vis absorption spectral changes for **5b** at 298 K ( $1 \times 10^{-4}$  M,  $\text{CH}_2\text{Cl}_2/\text{toluene} = 1/9$ , v/v). (c) Time-dependent UV/Vis absorption spectral changes for **5c** in  $\text{CH}_2\text{Cl}_2/\text{toluene}$  (1/9, v/v) at  $1 \times 10^{-4}$  M. (d) Time-dependent degree of aggregation for **5b** calculated from the absorbance at 480 nm at the indicated concentration in  $\text{CH}_2\text{Cl}_2/\text{toluene}$  (1/9, v/v).

UV/Vis absorption spectra for **5a** in the monomeric state (in  $\text{CH}_2\text{Cl}_2$ ) and aggregate state (in  $\text{CH}_2\text{Cl}_2/\text{toluene}$ , 1:9, v/v) are displayed in Figure 11. Complex **5a** in  $\text{CH}_2\text{Cl}_2$  solution ( $1 \times 10^{-4}$  M) shows two moderate absorption bands at 420

and 445 nm. Upon aggregation, a new broad absorption band at 485 nm emerged for **5a** and **5c**, while two new absorption bands at 450 and 480 nm appeared for **5b** (Figure 11b). Kinetically controlled supramolecular polymerization process were observed for **5b-c**, reflected by the time-dependent absorption spectral changes and the influence of concentration on the rate of the self-assembly process (Figure 11b-d). The kinetically controlled self-assembly of complex **5b** showed acceleration upon increasing the complex concentration (Figure 11d).

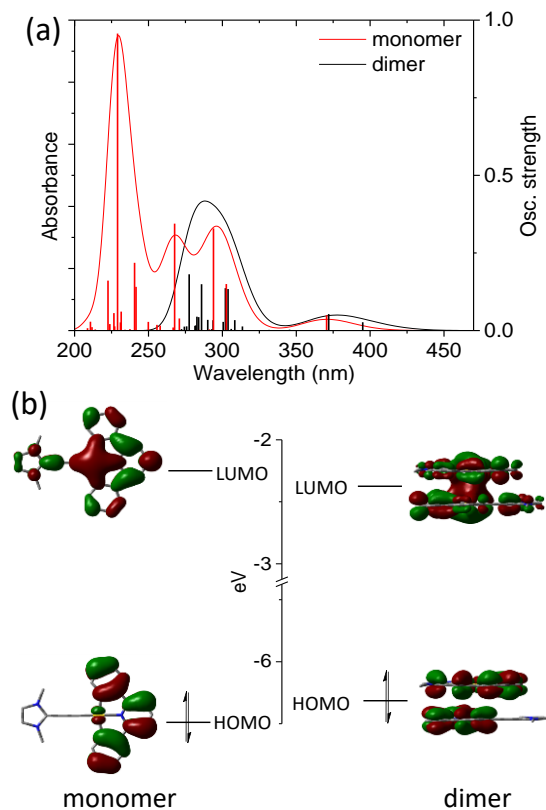
Complexes **5a-c** were found to be nearly non-emissive in the monomeric state, while they display moderate emission intensity in the aggregate state ( $1 \times 10^{-4}$  M,  $\text{CH}_2\text{Cl}_2/\text{toluene}$ , 1:9, v/v). The aggregates of **5a-c** exhibit a broad emission peak at 560, 605, and 565 nm, respectively, as shown in Figure 12. High phosphorescence quantum yields up to 61% have been reported for charge neutral pincer-type  $\text{Au}^{\text{III}}$  complex in solutions.<sup>49,50</sup> However, for cationic  $\text{Au}^{\text{III}}$  complexes, only moderate emission quantum yield of  $\sim 13\%$  could be achieved, such as that in  $[\text{Au}^{\text{III}}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})\text{L}]^+$  ( $\text{L} = \text{N}$ -heterocyclic carbene) system.<sup>51</sup> As the coordinated NHC ligand is orthogonal to the  $[\text{Au}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})]$  plane, intermolecular interaction and aggregation is disfavored. Thus the  $[\text{Au}^{\text{III}}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})\text{L}]^+$  system was not chosen for this work. According to the measured emission quantum yield (1%-6%) and excited lifetime (0.5-4  $\mu\text{s}$ ) for **5a-c**, the  $k_r$  (radiative decay rate) values were calculated to be  $1.5\text{-}2.1 \times 10^4 \text{ s}^{-1}$ , which are higher than those of reported  $\text{Au}(\text{C}^{\wedge}\text{N}^{\wedge}\text{C})$ -alkyl complexes (intraligand excited state in nature with  $k_r$  of approximately  $2 \times 10^3 \text{ s}^{-1}$ ).<sup>48</sup> The increased  $k_r$  values indicate the involvement of Au-metal characters in the frontier molecular orbitals of the excited state (LMMCT transition revealed by the DFT/TDDFT calculations discussed later), facilitating the ISC (inter-system crossing process) and the  $\text{T}_1 \rightarrow \text{S}_0$  transition.



**Figure 12.** (a) Emission spectra of **5a-c** in the aggregate state ( $1 \times 10^{-4}$  M,  $\text{CH}_2\text{Cl}_2/\text{toluene} = 1:9$ , v/v) at 298 K. Complexes **5a-c** in the monomeric state in  $\text{CH}_2\text{Cl}_2$  showed no observable emission. (b) From left to right: emission color of monomer-**5a** in  $\text{CH}_2\text{Cl}_2$  and aggregate-**5a**, aggregate-**5c**, and aggregate-**5b** in the mixed solvent ( $\text{CH}_2\text{Cl}_2/\text{toluene} = 1:9$ , v/v) at  $1 \times 10^{-4}$  M ( $\lambda_{\text{ex}} = 365$  nm). TEM images of (d) **5a**, (e) **5b** and (f) **5c** at  $1 \times 10^{-4}$  M in the mixed solvent of  $\text{CH}_2\text{Cl}_2$  and toluene (1:9, v/v).

## DFT/TDDFT Calculations and Excited States of Phosphorescent Au<sup>III</sup> Aggregates

DFT/TDDFT calculations were directly performed on the X-ray crystal structures of **1a** and **3d** without further optimization. As there is no available X-ray crystal structure for **4**, we first optimized the structures of **[4]<sub>2</sub>** and **[4]<sub>4</sub>** in the head-to-tail arrangement and then performed the subsequent calculations on the optimized structures. As revealed by the calculated absorption spectrum (Figure 13a) for **3**, the lowest absorption band for the dimer showed a slight redshift (from 371 nm for the monomer to 395 nm for the dimer). From experimental observations, the lowest absorption band of complex **3** redshifted from 400 to 420 nm upon aggregation. In the TDDFT calculations, the lowest energy absorption bands for the monomer and dimer are mainly from the HOMO→LUMO transitions. Further investigations by the molecular orbital (MO) diagrams (Figure 13b) of monomer **3** show the involvement of minor Au<sup>III</sup>(6p<sub>z</sub>) orbitals in the LUMO together with the major component of the  $\pi$  orbital from the ligand. In dimer **[3]<sub>2</sub>**, a weak bonding orbital  $\sigma$  (6p<sub>z</sub>) is proposed to be formed between two Au<sup>III</sup> centers and to become the new LUMO with lower energy. Therefore, the nature of the excited state of **3** upon aggregation is assigned to mixed minor <sup>1</sup>LMMCT and major <sup>1</sup>[ $\pi - \pi^*$ ] transitions. A similar case was found for complex **4** upon aggregation: The excited states for **[4]<sub>2</sub>** and **[4]<sub>4</sub>** originate from the mixed <sup>1</sup>LMMCT and <sup>1</sup>[ $\pi - \pi^*$ ] transitions (Figure S22). By contrast, for **1**, the excited state of the aggregated species is assigned to the <sup>1</sup>[LLCT] (ligand-to-ligand charge transfer) transition localized on the ligands with no Au<sup>III</sup> contribution (Figure S21).



**Figure 13.** Calculated (a) absorption spectra and (b) MO diagram of **3** and dimer **[3]<sub>2</sub>** in the S<sub>0</sub> state.

By adding electron-donating OMe groups to the pincer-type ligand of **4** to give **5**, the energy of  $\pi^*$  orbital (located at ligand C<sup>N</sup>^C) is destabilized, rendering more component of Au<sup>III</sup> 6p<sub>z</sub><sup>2</sup> orbital in the LUMO. Therefore, we observed the formation of a broad redshifted absorption/emission band (instead of a vibronic band) and increased *k<sub>r</sub>* constants at aggregate states for **5a**, which indicates the formation of an excited state with charge transfer properties and the involvement of more Au<sup>III</sup> 6p<sub>z</sub><sup>2</sup> orbital component in the LUMO.

## Discussion

### Au<sup>III</sup>-Au<sup>III</sup> Interaction in the S<sub>0</sub> and T<sub>1</sub> States and the Nature of the <sup>1</sup><sub>3</sub>LMMCT Transition

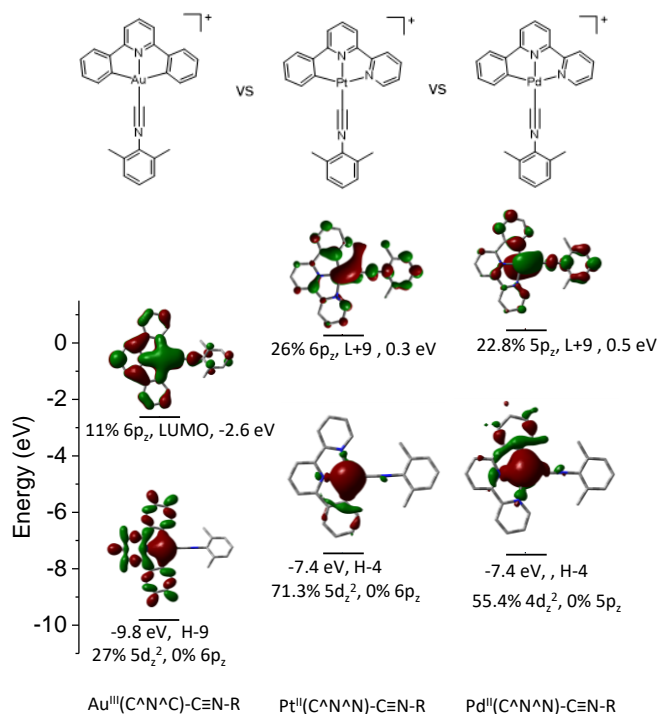
To address the question as to whether an attractive Au<sup>III</sup>-Au<sup>III</sup> bonding interaction exists or not in pincer-type Au<sup>III</sup> complexes, we also studied their isoelectronic counterparts, the pincer Pt<sup>II</sup> and Pd<sup>II</sup> complexes, which are well documented to display M-M interactions (in both the S<sub>0</sub> and T<sub>1</sub> states), for reference and comparison.

In some early semiempirical models,<sup>52</sup> the attractive closed-shell d<sup>8</sup>-d<sup>8</sup> interaction was considered to arise mainly from the orbital term<sup>53</sup> and was induced by the relativistic effects of heavy atoms, as reflected by the large radial contraction of s/p orbitals and radial expansion of d/f orbitals.<sup>54</sup> However, recent theoretical studies stated that M-M orbital interactions are less important than intermolecular dispersion interactions.<sup>8,53,55,56</sup> In this work, relativistic and dispersion effects on the Au<sup>III</sup>-Au<sup>III</sup> interaction in the S<sub>0</sub> ground state were considered and evaluated. By taking into account the relativistic effect of the heavy Au<sup>III</sup> atom, the 6p<sub>z</sub> orbital is stabilized and becomes lower in energy, compared to the one without considering the relativistic effect (Figure S19). Similar to recently reported computational studies, the calculation results indicate that the ligand-ligand dispersive interaction is the major driving force for the self-assembly process of Au<sup>III</sup> aggregates (Figure S20). Although closed-shell M-M orbital interaction is not always the driving force for the dimerization/oligomerization process of organometallic complexes, the interactions among different orbitals between two monomers will generate new electronic configurations in the dimer structure. Therefore, a distinct excited state with lower energy will appear for the dimer/oligomer species, which is one of the most important findings obtained by studying M-M orbital interactions. In the following section, to determine the electronic configurations and excited states in the dimer/aggregate forms of Au<sup>III</sup> complexes, main attention is paid to M-M orbital interactions (instead of M-M dispersion interactions or other noncovalent interactions).

For the d<sup>8</sup>-d<sup>8</sup> metal complexes where the metal is Pd<sup>II</sup> or Pt<sup>II</sup>, the overlap of two d<sub>z<sup>2</sup></sub> orbitals in the axial direction between two metal cations results in the destabilization of the orbital  $\sigma^*(d_{z^2})$ . Therefore, the antibonding  $\sigma^*(d_{z^2})$  between two metal cations becomes the new HOMO of the dimer/oligomer species.<sup>57</sup> A qualitative molecular orbital model for d<sup>8</sup>-d<sup>8</sup> metal complexes in the ground state is given in Figure S23.

However, the energy of the 5d<sub>z<sup>2</sup></sub>(Au<sup>III</sup>) orbital would be relatively lower for the Au<sup>III</sup> complexes studied in this work than for the Pt<sup>II</sup>/Pd<sup>II</sup> complexes, as the Au<sup>III</sup> cation is more electrophilic in nature. The HOMO of the pincer-type Au<sup>III</sup> complexes would mainly localize on the tridentate ligand. Weak d orbital overlap, if any, is expected in this series of Au<sup>III</sup> complexes. Moreover, the 6p<sub>z</sub> orbital of Au<sup>III</sup> would be

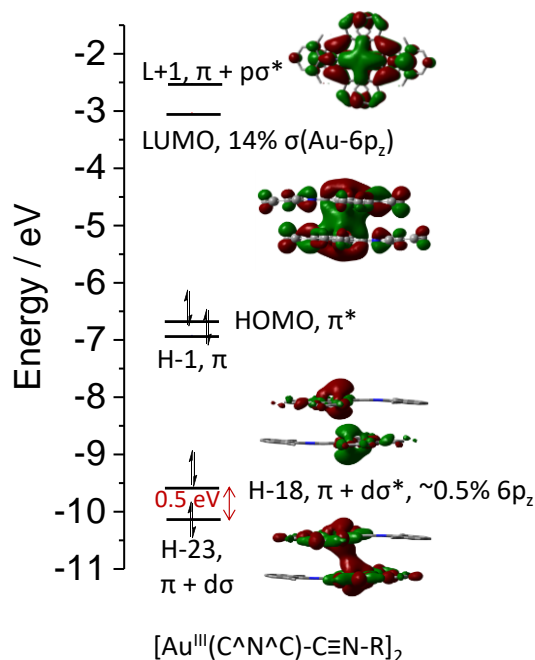
much lower lying in energy than the  $6p_z(\text{Pt}^{\text{II}})$  and  $5p_z(\text{Pd}^{\text{II}})$  orbital and is considered to make some contributions to the LUMO.



**Figure 14.** MO diagrams of **4** and its isoelectronic  $\text{Pt}^{\text{II}}$  and  $\text{Pd}^{\text{II}}$  counterparts  $[(\text{C}^{\wedge}\text{N}^{\wedge}\text{N})\text{MC}\equiv\text{N}(2,6\text{-dimethylphenyl})]$  ( $\text{M} = \text{Pt}$  or  $\text{Pd}$ ) in the optimized  $S_0$  structures.

We performed DFT calculations to compare  $\text{Au}^{\text{III}}$  complex **4** with its isoelectronic  $\text{Pt}^{\text{II}}/\text{Pd}^{\text{II}}$  counterparts  $[(\text{C}^{\wedge}\text{N}^{\wedge}\text{N})\text{MC}\equiv\text{N}(2,6\text{-dimethylphenyl})]$  ( $\text{M} = \text{Pt}$  or  $\text{Pd}$ ) in the ground state (Figure 14). The energies and proportions for  $nd_z^2$  and  $(n+1)p_z$  orbitals are displayed for all  $\text{Au}^{\text{III}}$  and  $\text{Pt}^{\text{II}}/\text{Pd}^{\text{II}}$  complexes. From Figure 14, the energy of the  $nd_z^2$  orbital in the  $\text{Au}^{\text{III}}$  complex is much lower than that of the  $\text{Pt}^{\text{II}}/\text{Pd}^{\text{II}}$  complexes (differs by approximately 2.4 eV), which is consistent with the electrophilic nature of the  $\text{Au}^{\text{III}}$  cation. For the  $(n+1)p_z$  orbital, the one in the  $\text{Au}^{\text{III}}$  complex is approximately 2.9 eV lower than that in its  $\text{Pt}^{\text{II}}/\text{Pd}^{\text{II}}$  counterparts, causing the  $6p_z$  orbital to become part of the LUMO in the  $\text{Au}^{\text{III}}$  complex.

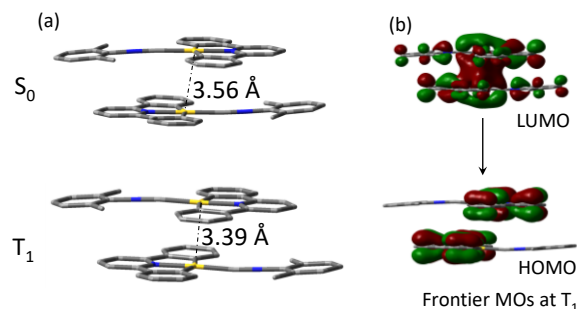
In the  $\text{Au}^{\text{III}}$  dimer's MO diagram (Figure 15), small splitting between the two  $\text{Au}(5d\sigma^*)$  and  $\text{Au}(5d\sigma)$  orbitals ( $\sim 0.5$  eV) emerged with a minor contribution from the  $\text{Au}(6p_z)$  orbital of approximately 0.5%, indicating a weak Au-Au orbital interaction in the ground state (compared to the Pt-Pt/Pd-Pd orbital interactions shown in Figure S25). At the same time, the bonding  $6p\sigma$  orbital became the new LUMO in the  $\text{Au}^{\text{III}}$  dimer, and therefore, a new excited state will be generated upon excitation. A qualitative electronic configuration of the  $\text{Pt}^{\text{II}}\text{-Pt}^{\text{II}}/\text{Pd}^{\text{II}}\text{-Pd}^{\text{II}}$  and  $\text{Au}^{\text{III}}\text{-Au}^{\text{III}}$  dimers in excited state  $T_1$  is displayed in Figure S24. In the case of the  $\text{Au}^{\text{III}}$  dimer, upon excitation, an electron in the HOMO or a lower energy orbital (ligand  $\pi$ ) would be excited to the LUMO, which can be regarded as a weak bonding orbital between two  $\text{Au}^{\text{III}}$  centers. Similar to the  $\text{Pd}^{\text{II}}/\text{Pt}^{\text{II}}$  dimers, a half bond (or less, depending on the metal character in the LUMO) was proposed to be formed in this  $\text{Au}^{\text{III}}\text{-Au}^{\text{III}}$  system in the  $T_1$  excited state.



**Figure 15.** MO diagram of  $\text{Au}^{\text{III}}$  dimer  $[\mathbf{4}]_2$  in the optimized  $S_0$  structure.

The attractive  $\text{Au}^{\text{III}}\text{-Au}^{\text{III}}$  bonding interaction in the  $T_1$  excited state was further elucidated by the DFT/TDDFT calculation results. The optimized triplet excited state geometry of  $[\mathbf{4}]_2$  reveals a shortening of the Au-Au distance from 3.56 Å in  $S_0$  to 3.39 Å in the  $T_1$  state (Figure 16). This finding is supportive of the existence of attractive  $\text{Au}^{\text{III}}\text{-Au}^{\text{III}}$  interactions in the  $T_1$  state.

The nature of the emissive excited state for the  $\text{Au}^{\text{III}}$  dimer in  $T_1$  was assigned to the mixture of major  ${}^3[\pi - \pi^*]$  and minor  ${}^3\text{LMMCT}$  transitions shown in Figure 16b. From the MO diagrams, it can be noticed that the  $6p\sigma$  orbital delocalized onto the two  $\text{Au}^{\text{III}}$  centers in the LUMO. The emission energy of  $[\mathbf{4}]_2$  in the  $T_1$  state was calculated to be 2.3 eV, which is close to the experimental value of 2.0 eV (Figure 9).



**Figure 16.** (a) Optimized dimer structures of  $[\mathbf{4}]_2$  in the  $S_0$  (top) and  $T_1$  (bottom) states. (b) FMO diagram of  $[\mathbf{4}]_2$  in the optimized  $T_1$  structure.

Although a bonding metal-metal interaction was considered formed in both of the  $\text{Pt}^{\text{II}}\text{-Pt}^{\text{II}}/\text{Pd}^{\text{II}}\text{-Pd}^{\text{II}}$  and  $\text{Au}^{\text{III}}\text{-Au}^{\text{III}}$  cases in the  $T_1$  state, the nature of the  ${}^3\text{MMLCT}$  and  ${}^3\text{LMMCT}$  excited states is distinctly different. In the  $\text{Pd}^{\text{II}}/\text{Pt}^{\text{II}}$  case, upon excitation, the  $\text{Pd}^{\text{II}}/\text{Pt}^{\text{II}}$  cation would be oxidized, while the  $\pi$ -conjugated ligand would be reduced as the



charge transfers from the metal to the ligand (MMLCT). By contrast, for the Au<sup>III</sup>-Au<sup>III</sup> dimer in the excited state, the metal cation would be reduced by the ligand during the reverse charge transfer flow of the LMMCT. The change in the charge transfer flow was considered to have significant influence on the photophysical, photochemical and semiconducting properties of the d<sup>8</sup> organometallic complexes. We conceive that Au<sup>III</sup> complexes possess promising potential in the design of supramolecular materials with new functions as a result of their unique electronic properties and distinct excited states.

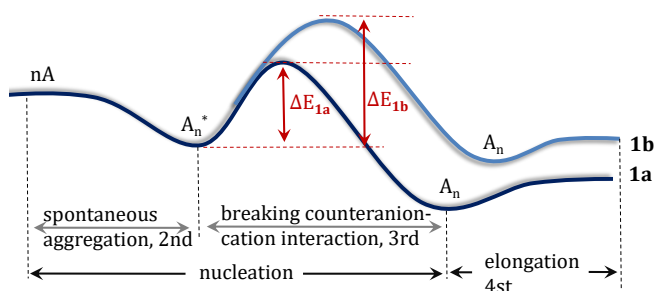
### Pathway Complexity of Complexes 1 and 3

To account for the counteranion- and/or concentration-dependent kinetics and supramolecular polymer morphologies observed for **1** and **3**, a qualitative kinetic model is applied.<sup>28,43</sup> As shown in Figure 17, symbol A represents the monomer, A<sub>n</sub> is the aggregate species having n monomers, and n represents the size of the nucleus.

A four-step self-assembly mechanism was used to account for the systems studied here: (Step 1) At first, Au<sup>III</sup>-counteranion complex is soluble in CH<sub>3</sub>CN/CH<sub>2</sub>Cl<sub>2</sub> and would be surrounded by solvent molecules in these solvents. The counteranion has relatively close contact with complex cation in the organic solvent cage at this stage. (Step 2) Upon injecting a poor solvent (water or toluene) into a solution (CH<sub>3</sub>CN or CH<sub>2</sub>Cl<sub>2</sub>) of the complex, the solvent cage is disrupted and the monomers would be quickly trapped in the kinetic aggregate state A<sub>n</sub><sup>\*</sup>, termed the “inactive nucleus”. As there is no decent energy barrier between monomer A and inactive nucleus A<sub>n</sub><sup>\*</sup>, this process occurs very quickly and spontaneously. The counteranion still has relatively close contact with Au<sup>III</sup> complex at this kinetic aggregate state. (Step 3) By breaking the electrostatic interaction between the counteranion and complex cation, the inactive nucleus is slowly transformed into “active nucleus” A<sub>n</sub>. This process is counteranion dependent, and the activation energy barrier is significantly affected by the interaction between the complex cation and counteranion. This transformation process is also considered to be the rate-determining step which is responsible for the induction period (lag time) observed for **1b** and **3a** (Figure 1d & 6b). (Step 4) If there are excess remaining monomers in the solution, these monomers will further aggregate at the terminals of the activated nucleus, which is termed the elongation process as shown in Figure 17.

In summary, there are three types of molecules and four types of intermolecular interactions involved during the kinetically controlled self-assembly process: cationic Au<sup>III</sup> complex, counteranion and solvent molecule (mixed solvent); cation-cation interaction, cation-counteranion interaction, counteranion-counteranion interaction and solvation energy (which could be further decomposed to: kinetic aggregate’s solvation energy, thermodynamic aggregate’s solvation energy and monomer’s solvation energy). The transformation from kinetic aggregate to thermodynamic aggregate is accelerated by the following non-covalent interactions: strong cation-cation dispersive interaction; weak cation-anion interaction; less kinetic-aggregate’s solvation energy and more thermodynamic-aggregate’s solvation energy. Change of the metal complex (ligand or metal) influences the cation-cation interactions (Ligand-Ligand dispersive, Metal-Ligand dispersive and M-M closed shell interaction); change of the counteranion influences the cation-counteranion interaction.

Change of the solvent (ACN/water or DCM/toluene) affects the solvation energy.



**Figure 17.** Qualitative energy landscape of the supramolecular polymerization models for **1a** and **1b**.

The following equation, developed by Oosawa and Kasai, was employed to describe the reaction rate of nucleus formation (step 3):<sup>43,58,59</sup>

$$\frac{d[A_n]}{dt} = k^*[A]^n - k_-^*[A_n] - k[A][A_n] + k_-[A_{n+1}] \quad (1)$$

where [A] is the concentration of the monomer; [A<sub>n</sub>] is the concentration of the active nucleus; k<sup>\*</sup> and k<sub>-</sub><sup>\*</sup> represent the rate constants for the formation and dissociation of A<sub>n</sub> (step 3), respectively; and k and k<sub>-</sub> are the rate constants of the forward and reverse chain elongation reactions (step 4), respectively.

For the formation rate of the activated nucleus and elongated polymer (i ≥ n), we have:

$$\frac{dc_p^*}{dt} = \frac{d \sum_{i=n}^{\infty} [A_i]}{dt} = k^*[A]^n - k_-^*[A_n] \quad (2)$$

where c<sub>p</sub><sup>\*</sup> is the concentration of all polymers equal to or larger than n.

**Concentration-Dependent Kinetics for 1a and 3a.** According to equations (1) and (2), the formation rates of the active nucleus and elongated polymer are positively correlated with the total monomer concentration [A]. At a high concentration of **1a** and **3a** (1 × 10<sup>-4</sup>M), as the nucleation rate is very fast, the monomers will be quickly consumed and transformed into the active nucleus. The depolymerization of the nucleus and elongated polymers can be approximately ignored in this situation. All of the complexes aggregating at the terminal of the nucleus will be instantly activated. Therefore, only short, uniform 1D nanowires emerged (Figure 2a and Figure 7a).

However, upon lowering the concentration of **1a** and **3a**, the rate of the nucleation process (step 3) was slowed down, indicating that excess monomers may remain in the system, together with the transformed active nucleus A<sub>n</sub>. The depolymerization of the nucleus cannot be ignored in this situation. Only part of the complexes (at the end of the nucleus) can be activated at a relatively slow rate. Gradually, the remaining monomers in the solution began to elongate along only the activated part of the nucleus (step 4); therefore, 2D supramolecular structures were observed for **1a** and **3a** at relatively low concentrations.

**Counteranion-Dependent Kinetics for 1.** According to our previous work,<sup>16, 22</sup> PF<sub>6</sub><sup>-</sup> (compared with ClO<sub>4</sub><sup>-</sup>) would show weaker interactions with the metal complex, leading to a smaller kinetic energy barrier (ΔE<sub>1a</sub>) for **1a** (counteranion: PF<sub>6</sub><sup>-</sup>) than for **1b** (counteranion: ClO<sub>4</sub><sup>-</sup>).

According to the Arrhenius equation, the rate constant,  $k$ , for a chemical reaction is inversely proportional to the activation energy,  $\Delta E$ . Therefore, we have:

$$k^*_{1b} < k^*_{1a} \quad (3)$$

According to equation (2), the reaction rate of the polymerization is positively related to the forward rate constant  $k^*$ . When  $k^*_{1b} < k^*_{1a}$ , the nucleation process (step 3) for **1b** will be seriously retarded, which was reflected by the long lag time shown in Figure 1d. During the lag time, a large number of complexes would remain as monomers in the system because of the dissociation of the inactive nucleus. Similarly, only part of the complexes (at the end of the nucleus) can slowly be activated. Upon elongation of the remaining monomers, the 2D bowknot-shaped nanostructure shown in Figure 3b was formed and observed.

## CONCLUSIONS

We demonstrated the kinetically controlled self-assembly of phosphorescent pincer-type  $d^8$  Au<sup>III</sup> aggregates by tuning the counteranion and concentration. Living supramolecular polymerization was achieved for Au<sup>III</sup> complexes with acetylene as the ligand. A qualitative kinetic model was applied to explain the pathway complexity of Au<sup>III</sup> complexes studied in this work. A close Au-Au contact of 3.367 Å was observed in the X-ray crystal structure of the Au<sup>III</sup>-allenylidene complex. A DFT/TDDFT study revealed that pincer Au<sup>III</sup>-isocyanide and Au<sup>III</sup>-allenylidene complexes can display weak attractive Au<sup>III</sup>-Au<sup>III</sup> interactions upon excitation; furthermore, a mixed  $^{1/3}$ LMMCT and  $^{1/3}$   $\pi$ - $\pi^*$  excited state was assigned to the polymeric Au<sup>III</sup> species through the supramolecular polymerization.

## ASSOCIATED CONTENT

### Supporting Information

Experiment procedures, characterization of complexes (including NMR spectra, high resolution mass data and element analysis), computational and experimental details, molecular coordinates by calculations, emission and excitation spectra, UV/Vis absorption spectra, TEM images, and CIF file for the X-ray crystal structure of **3d**.

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### Notes

The authors declare no competing financial interest.

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