

1 **Self-regulated symmetric breaking enables full-space phototaxis of animate  
2 materials**

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18

19 **Abstract**

20 Living organisms in their early forms adopted simple structures and could adapt to different cues  
21 from the ambient environment in a highly regulated and efficient manner. However, synthetic  
22 systems have achieved less with more human intervention than their natural counterparts. Here,  
23 we show that a simple animate material can recognize the directions of constant photonic  
24 illumination and follow a path with strong directional regulation inside a constraint-free, fluidic  
25 space. By spontaneously breaking the symmetry of multiple surrounding energy fields, the  
26 nanostructured stimuli-responsive polymers-based animate material swiftly assumes an optimal  
27 pose and creates directional flow around itself, which it follows to achieve robust full-space  
28 phototaxis. In addition, this phototaxis enables a series of complex, self-adaptive locomotion  
29 underwater. We demonstrate that this versatility is empowered by the fast reconFigurability and  
30 synergy of photo-fluidic interactions composing closed-loop self-control. The untethered,  
31 ambient-powered animate-material manoeuvres through obstacles agilely, following the cues of  
32 constant illumination without any electronics or close-range human intervention.

33

34 Planktonic microbes move hundreds of metres towards the surface each night to feed before  
35 retreating back down at dawn<sup>1,2</sup>. Many of them have developed receptors for directed locomotion,  
36 known as phototaxis, gyrotaxis, rheotaxis, and chemotaxis, so they can sense many ambient  
37 gradients and move along a preferred direction to acquire food and the chance for reproduction  
38 and move away from danger in an adaptive and sustainable manner<sup>3,4</sup>. This behaviour links the  
39 epipelagic and mesopelagic zones, critically contributing to the oceanic biological carbon pump<sup>5,6</sup>.  
40 Inspired by nature, recent years have witnessed extensive pioneering efforts to advance robots in a  
41 fluidic environment across a great range of scales<sup>7-14</sup>. However, the controls of the artificial  
42 systems still rely on preprogramed electronics and a high level of human attention at close range,

43 *i.e.*, magnetic field manipulation<sup>15,16</sup>, accurate aiming<sup>10,17</sup>, flash frequencies of lasers<sup>8,18</sup>, and preset  
44 constraints in space<sup>17,19</sup>. These prerequisites and special conditions do not exist in nature and  
45 hence set obstacles to practical application.

46 Breakthroughs have been achieved in self-regulated materials<sup>20,21</sup> for autonomous phototropic  
47 tracking<sup>22</sup> and oscillation<sup>17</sup>. The demonstrated photophobic taxis was constrained to a water  
48 surface and within a one-dimensional fluidic channel<sup>17</sup>. By far, synthetic systems that are capable  
49 of omnidirectional, full-space phototaxis without electronics or human intervention remain a  
50 highly desirable goal with the aim to design an embodied cognitive material<sup>23-25</sup> in free space  
51 rather than a reflexive material under many limitations.

52 Here, we develop a monolithic, animate-material system that recognizes and moves in any  
53 direction of constant photonic illumination (**Fig. 1**). As shown in **movie S1**, the untethered soft  
54 material identifies the direction of irradiation and move towards it, submerged in an aqueous  
55 environment. While illuminated, the soft robot harnesses the photonic energy and breaks the  
56 symmetry of temperature, flow fields, and morphology. The multi-physical symmetric-breaking  
57 allows the animate material to sense the direction of the incidence (**Fig. 1a**), create a phototropic  
58 stream that provides propulsion (**Fig. 1b**), and spontaneously assume the optimal posture and gait  
59 (**Fig. 1c**) for swift phototactic locomotion.

60 During the symmetric-breaking-induced locomotion, the animate material is capable of  
61 self-regulating the photonic energy-input by adaptively enhancing convective heat loss around  
62 itself, which provides negative feedback for the correction of deviated traveling direction (**Fig. 1d**,  
63 **e**) to implement robust phototactic manoeuvres. The material system employs an  
64 ultra-high-sensitivity stimuli-responsive material that recovers at an unprecedented rate to allow a  
65 fast response to any new event (**Fig. 1f**), offering a steering finesse of full-space phototaxis that is  
66 completely based on the cognition of the material itself (**Fig. 1g**).

67

#### 68 **Spontaneous symmetric breaking and recovery**

69 We adopt a simple geometry similar to a jellyfish with axial symmetry that comprises a  
70 hemispherical bell and six vertical tentacles. The monolithic structure of the jellyfish-like  
71 phototactic vehicle (JPV) permits facile additive manufacturing technology, such as P $\mu$ SL, as well  
72 as moulding manufacturing for mass production (**Supplementary Fig. S3**). The bell contains an  
73 air bubble as a key component for balancing gravity and breaking the symmetry of the density  
74 distribution (**Fig. 1c**). Without directional illumination, the bubble is kept in an up-centre position,  
75 providing buoyancy to partially balance gravity and keeping the JPV in an upright position.

76 The universal design of JPVs can be fulfilled with many active materials, *i.e.*, light-active  
77 polymers<sup>26-28</sup> and electroactive polymers<sup>29-32</sup>. Here, we select the stimuli-responsive  
78 poly(N-isopropylacrylamide) (PNIPAAm)-based hydrogel<sup>33</sup> as an example for its iconic,  
79 reconFigurable morphing across its low critical solution temperature (LCST). Similar to the  
80 sensory system spread all over the body of a squid<sup>34</sup>, the hydrogel is homogeneously blended with  
81 highly efficient photo-thermal nanoparticles, such as Au nanoparticles (AuNPs) and reduced  
82 graphene oxides (r-GO), to ensure omnidirectional sensing of incident photons<sup>35</sup> (**Supplementary**  
83 **information Section 1**).

84 Once illuminated, three symmetries of the JPV are broken simultaneously: 1. a temperature  
85 gradient is established along the direction of the illumination, due to the effective photothermal  
86 coupling that heats the illuminated surface (**Fig. 2a** and **Supplementary Fig. S7**); 2. a convective

87 flow field is generated from the heated surface (**Fig. 2b**); 3. when the local temperature is higher  
88 than the LCST, the local shrinkage of the gel squeezes the embedded bubble balancer away from  
89 its symmetric position in the bell (**Fig 1c, 2c** and **Supplementary Video 2**), introducing a  
90 redistribution of density and hence a phototropic tilting towards the light source (**Fig. 2d**). The  
91 phototropic behaviour was not observed in the control sample made of thermally inactive  
92 poly(acrylamide) (**Supplementary Fig. S12**) with the same amount of photo absorbers. We  
93 employ a multi-physics model based on the finite element method (FEM) to investigate the  
94 synergy of the multi-energy coupling that governs symmetric breakings (**Supplementary Fig. S7**).  
95 The mechanical strain and temperature gradient showed great agreement with the experimental  
96 results (**Fig. 2a, d**). As shown in **Fig. 2d**, the off-centred bubble leads to a rapid rearrangement of  
97 the JPV's centre of buoyancy (B), causing the metacentre (M) to move below the centre of gravity  
98 (O), generating a clockwise torque under illumination from the right. The torque exhibits a  
99 self-adaptive lean towards the direction of illumination, offering reduced fluidic drag that favours  
100 phototaxis (**Supplementary Fig. S20c**), as the JPV identifies the direction of the incidence and  
101 prepares itself for movement. Our theoretical analysis shows that the drag coefficient is reduced  
102 by 30% when the tilting angle changes from 90° (upright position) to 50° (**Supplementary Fig.**  
103 **S20**).

104 Agile phototactic behaviour demands ultrafast recovery to prevent a time lag resulting from a  
105 varying direction of illumination (**Fig. 2e**). However, the re-swelling rate of PNIPAAm-based  
106 hydrogels is notoriously slower than their shrinking rate. To restore the chronic symmetry of  
107 actuation and recovery, we modify the porous size and porosity of the gel by tightly controlling  
108 the chemical and physical processes of fabrication. For quantitative analysis, we tested the rate of  
109 phototropic actuation and recovery of the tentacle for different angles of incidence. At an angle of  
110 20°, the ice-templated, hybrid-crosslinked hydrogel pillar fully recovers less than a second,  
111 which is 30 times faster than the best previously reported values<sup>34</sup>, under similar pre-conditions  
112 (**Fig. 2f, g, and Supplementary Information Section 2**).

113

#### 114 **Photothermal-tactic flow**

115 Planktonic organisms, ranging from bacteria to jellyfish, move around by following the flows,  
116 waving or rotating their flagella or lappet, mostly at a low Reynolds number. The JPV is designed  
117 to harness the momentum from the ambient light that translate into photothermal-tactic flows.  
118 Their tentacles provide functions such as balance keeping and grabbing, rather than providing  
119 propulsion. In nature, propulsion efficiency at low Reynolds numbers is low (~1%–3%)<sup>36</sup>,  
120 fundamentally because the thrust is generated by the same mechanism that resists motion: viscous  
121 drag. On the other hand, plankton are capable of moving hundreds of metres by adaptively  
122 positioning themselves in a favourable current. Hence, generating photothermal tactic flow around  
123 the JPV (mostly a passive element) is an important task for achieving long-distance, underwater  
124 phototaxis.

125 In the JPV, the momentum is originated from the temperature gradient, which defines the  
126 direction of the photothermal flow. Maintaining the symmetric-broken temperature field requires  
127 synergy between the photo-thermal environment inside and around the JPV, *e.g.*, the thermal  
128 conductivity, photonic penetration length, convection, and characteristic size, *etc*. Therefore, we  
129 studied Biot number of the JPV. As most plankons, JPV contains majorly water holding inside  
130 polymer chains. The skin of the JPV exhibits high absorption hence the illuminated surface can be

131 treated as a surface heat source. We find that the photothermal-tactic flow induced acceleration of  
132 the JPV is scaling insensitive (**Supplementary Fig. S22**) as long as the JPV is larger than a critical  
133 size (1 mm in current design) that is capable of maintaining effective symmetry breaking  
134 (**Supplementary Fig. S24**). The observation is consistent with the analysis of size-dependent Biot  
135 number ( $Bi$ ). Further reducing thermal conductivity of the JPV body could extend the critical size  
136 down to the scale of tens of micrometres (**Fig. 2h**).

137 As schematically shown in **Fig. 2i**, when illuminated from the above, the JPV quickly ascends  
138 to the surface (**Fig. 2j**), reminding us of the vertical migration of marine plankton that represents  
139 the largest biomass transport on Earth<sup>37</sup>. Homogeneously distributed photo-thermal nanoparticles  
140 efficiently transduce photonic energy from all directions to thermal energy<sup>38</sup>. The elevated  
141 temperature on the illuminated surface creates a convective flow indicating directional  
142 information (**Fig. 2k**). To examine the hydrodynamic details, we conducted particle image  
143 velocimetry (PIV) to monitor the flow field that carries the JPV. The results are in good agreement  
144 with FEM-based computational fluid dynamics (CFD) simulation (**Fig. 2k, l**).

145 The PIV images of the hydrodynamic footprint reveal that the velocity distribution of the  
146 surrounding fluid moves as the JPV moves (**Fig. 2m and Supplementary Video 3**). We observe  
147 vortices during the phototactic movement of the JPV, which is considered a sign of inertial  
148 swimming<sup>10</sup>. The convective flow occurring around the JPV can be recognized as laminar flow  
149 based on the Reynolds number  $Re=uD/v$  (where  $u$  refers to the gliding velocity,  $D$  is the  
150 characteristic size, and  $v$  is the kinematic viscosity of the water). The PIV data indicate that the  
151 JPV is swimming in the low-Reynolds-number region where its natural counterparts live. During  
152 vertical ascendance, the submerged JPV is subjected to gravitational force  $G$ , buoyancy force  $F_B$ ,  
153 drag force  $F_D$ , and lift force  $F_L$ , and this upwards locomotion away from the bottom occurs as long  
154 as  $F_B+F_L>F_D+G$  (**Supplementary Fig. S22**). Vertical phototaxis is intensity-cognitive  
155 (**Supplementary Video 4 and 5**). The JPV can sense the power density of the photo-illumination,  
156 lean, and steer towards the brighter zone when the light field is not evenly distributed in space  
157 (**Supplementary Fig. S15 and Video 5 demo1**).

158 The intensity cognition is attributed to the synergy of multiple energy interactions, including  
159 photothermal energy transduction, thermochemical phase transition, chemo-mechanical  
160 deformation, thermofluidic convection flow, and flow-soft matter interaction. The synergy allows  
161 the JPV to maintain exposure and lean towards the brightest irradiation, which results in a high  
162 velocity of phototactic movement (**Supplementary Fig. S14a-d**). Constrained by the size of the  
163 container, the highest speed of the JPV is 18 mm/s and 1.5 BLs/s (**Supplementary Fig. S14d and**  
164 **Video 4**), which is significantly higher than those of many electric powered and tethered  
165 underwater robotics (**Supplementary Fig. S14e and table S1**) and comparable to those of many  
166 mammalian and non-mammalian organisms in nature that typically move at speeds exceeding 1  
167 BL/s<sup>39</sup>.

168 Compared to vertical phototaxis, phototaxis along any other direction (**Fig. 2n**) is more  
169 complicated to achieve due to temperature-induced upwelling. Under horizontal irradiation from  
170 the right, the symmetry of the temperature and flow field are broken, producing higher  
171 temperature and fluid velocity on the illuminated side (**Fig. 2o**). The temperature gradient near the  
172 illuminated surface reflects the directional information of the incidence but is also coupled with  
173 the adjacent flow field (**Fig. 2p**). The CFD and PIV results (**Fig. 2q, r**) indicate that the flow field  
174 consists of a significant part of the upwards convective flow that carries the JPV upwards and to

175 the right (**Fig. 2r**), resulting in inaccurate horizontal phototaxis. Additional measures should be  
176 taken for the JPV to self-regulate the upward flow, *e.g.*, reducing its surface temperature  
177 on-demand as negative feedback<sup>20</sup>, in order to continuously move in the direction of illumination.

178 **Directional self-regulation**

179 As the major provider of the fluidic thrust, the temperature gradient near the illuminated surface is  
180 defined by the input photonic heating power and the cooling power output to the ambient. We  
181 present two mechanisms to regulate the local temperature, hence enabling the locomotion of the  
182 JPV by allowing the JPV to self-manipulate the ambient convective heat transfer during  
183 movement.

184 Successful horizontal phototaxis is achieved, as shown in **Fig. 3a** and **Supplementary Video 1**.  
185 Under 532-nm laser irradiation, the AuNP-incorporated JPV automatically leaps up and forwards  
186 and falls back to the bottom. During each jumping cycle, only the horizontal displacement  
187 accumulates via phototaxis. The hopping cycle can be initiated under irradiation lower than 1000  
188 W/m<sup>2</sup> (**Supplementary Fig. S28a, b**). To further investigate the self-regulatory behaviour, we  
189 simulated and monitored the speed and the surface temperature of the JPV *in situ* during the  
190 phototactic hopping cycles. The *in-situ* monitored correlation between the speed and the local  
191 temperature of the JPV is in good agreement with the CFD simulation (**Supplementary Fig. S28c,**  
192 **d**). The speed and temperature-correlated convective heat transfer coefficient can be obtained as  
193 the mixed Nusselt number ( $Nu=(Nu_n^4+Nu_f^4)^{1/4}$ )<sup>40,41</sup>, where  $Nu_n$  is the natural convection and is  
194 related to the Rayleigh number ( $Ra$ ), and  $Nu_f$  refers to forced convection that is determined by the  
195 Reynolds number ( $Re$ ) (**Supplementary Information Section 3.2**).

196 To study the multi-physical, photo-thermo-fluidic dynamics, we show the time evolutions of  
197 scalar quantity of speed ( $v$ ), local temperature ( $T$ ), convective heat transfer coefficient ( $h$ ), input  
198 photonic power ( $P$ ) and convective cooling power ( $q$ ) near the illuminated bell area for two  
199 consecutive hopping cycles. When the JPV is ascending and subsequently moving away from the  
200 laser, the temperature saturates, and the speed of the JPV drops to zero, indicating strong photonic  
201 regulation. After a short period of inertial ascendance, the JPV falls back and completes the  
202 negative feedback cycle. The self-avoidance of light incidence indicates a strong photo-thermal  
203 regulation, which is quite similar to the characteristic thermal dance of shovel-snouted lizards  
204 (*Meroles anchietae*) to avoid overheating of their feet.

205 Notably, the ascending JPV exposes its photo-responsive tentacles to the laser beam and aligns  
206 with the incidence automatically (**Fig. 3a, b**) as the JPV controls its direction of movement. The  
207 bent tentacle strengthens the leaning pose, which helps protect against upwards flow, facilitate  
208 horizontal motion.

209 Compared to the narrow laser beam, which requires constant human involvement to keep  
210 aiming at the JPV, phototaxis towards a wide beam-width incidence may represent a more general  
211 case in the natural environment. The challenge is to achieve thermal regulation when the JPV can  
212 no longer hide from the light. Surprisingly, we observe similar hopping cycles (**Fig. 3c**) of the JPV  
213 illuminated by a spotlight. The speed and thermal analysis (**Fig. 3d**) show that the JPV is diving at  
214 a faster speed (2.5-fold) than that corresponding to laser illumination with the same total energy  
215 (**Supplementary Fig. S9**). The enhanced speed can be attributed to the continuous illumination of  
216 the spotlight that accumulates more momentum than in the case of self-avoiding-induced  
217 intermittent illumination. The high-speed JPV encounters ambient cold fluid more rapidly, in **Fig.**  
218 **3d**, we show that the time-dependent  $h$  is strongly correlated with the time evolution of the

real-time speed, suggesting speed-induced thermal regulation. Nusselt number corresponding to the maximum velocity of photonic and thermal regulation are comparable, however, the Richardson number ( $Ri=Gr/Re^2$ , representing the importance of natural convection relative to forced convection, where  $Gr$  is Grashof number and  $Re$  is Reynolds number) of photonic regulation is more than 20 times higher than that of thermal regulation, indicating that they are dominated by natural and joint convection, respectively (**Fig. 3e**). The much-increased convective heat transfer ( $hA(T_s - T_f)$ ) regulates the surface temperature and fulfils a negative feedback loop (**Fig. 3f**). By utilizing effective thermal regulation, full-space and omnidirectional phototaxis is achieved for any 3-D incident angle (**Supplementary Fig. S29 and Video 1 demo2**).

In addition, the JPV offers various hopping frequencies for selection by manipulating its excess gravity against the input photonic power. By reducing the volume of the bubble in the bell, the hopping frequency continues to decrease until an extreme case is reached in which the JPV slides horizontally instead of hopping (**Supplementary Fig. S31 and Video 6**). The tuning of the excess gravity is an effective tool to preset the phototactic behaviour according to the different ambient light, thermal and fluidic conditions in practical field deployment.

234

### 235 **Demonstration of complex locomotion**

236 The highly regulated phototaxis of the JPV can be characterized by the distance of the accurate  
237 tracking and the angular resolution. Distinguished from acquiring momentum from paddling a  
238 single tentacle that can be easily misaligned with the laser beam, the JPV performs robust  
239 phototaxis over a long distance without directional deviation (**Fig. 4a, b and Supplementary Fig.**  
240 **S33**). Limited by the dimension of the experimental space, the JPV exhibits accurate phototaxis of  
241 300 mm compared to 4.3 mm for single-tentacle paddling. In addition, by carefully cascading the  
242 strokes by 6 off-centre-positioned tentacles, the JPV adapts to exhibit rotational advance (**Fig. 4c**  
243 **and Supplementary Video 7**). The bent tentacle counteracts the overturning moment and hence  
244 provides additional balance when the JPV is climbing a ramp, further ensuring the stability of the  
245 movement (**Fig. 4d and Supplementary Video 8**). These tentacles can also serve as photonic or  
246 thermal sensitive grabbers that can track and grasp a bright or hot object (**Supplementary Fig.**  
247 **S13**).

248 Due to the spontaneous symmetric breaking and fast recovery of the polymeric matrix (**Fig. 1g**),  
249 the directional response of the JPV is instant. We observe that the JPV is capable of agile  
250 manoeuvring with fine directional and positional resolution. Following a constant light source  
251 with continuously varying angles, the JPV travels perfectly along with the complex strokes of  
252 Chinese calligraphic characters “自然” (**Fig. 4e and Supplementary Video 9**), which means  
253 “nature” and is written in the cursive script style (with connected characters). In addition to the  
254 capability of fine steering, we further demonstrate long-distance operation of the JPV with a  
255 beam-expanded laser placed 10 m away (**Supplementary Fig. S18**).

256 Different from the phototropic taxis when submerged, the JPV shies away from the light once it  
257 is floating on the water-air interface (**Fig. 4f**). The photophobic behaviour is induced by  
258 Marangoni convection, which is attributed to the symmetric breaking of surface tension on the bell  
259 at the water surface. Both phototactic modes are capable of accurately following the paths of the  
260 character “SJTU” (**Supplementary Fig. S32c, d**)

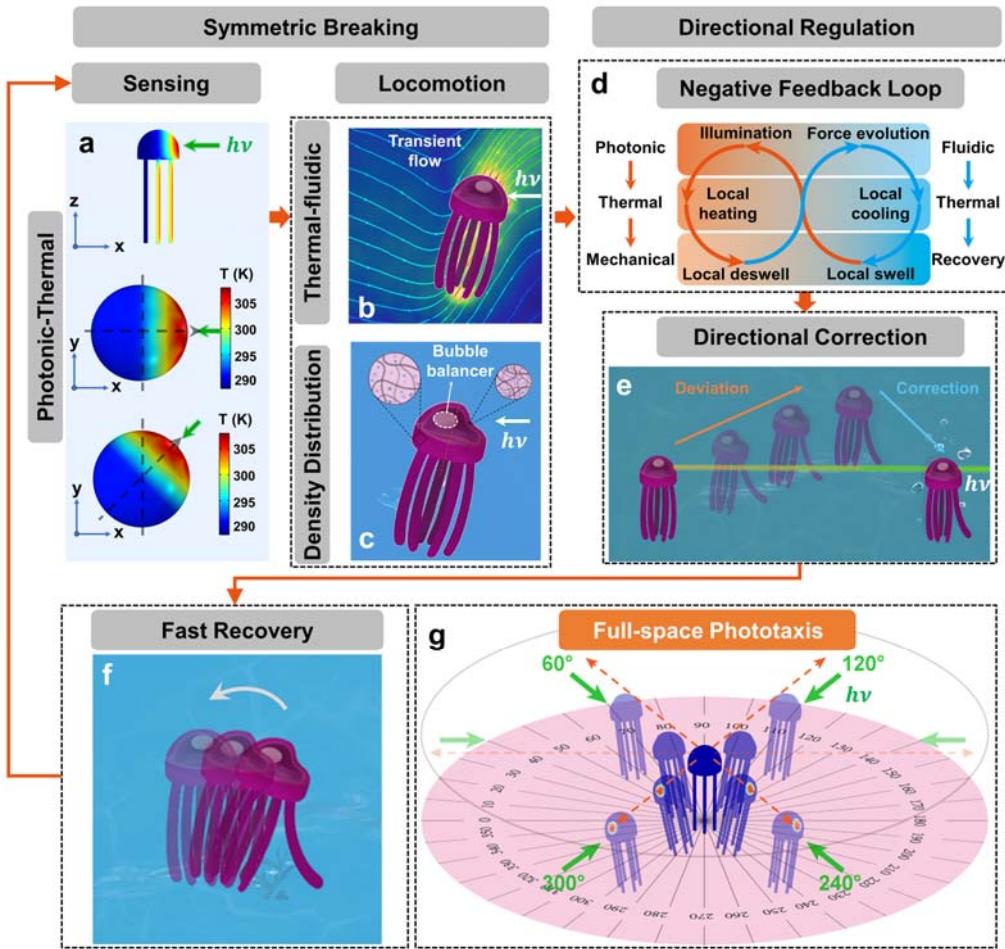
261 Combining the two phototactic behaviours, the JPV can dive up to the water surface and move

262 away from the light to avoid obstacles (**Supplementary Fig. S33 and Video 10**) and then dive  
263 back in and move towards the light (**Supplementary Video 11**). The cyclic behaviour is similar to  
264 the diurnal vertical migration and cycling of many planktonic biomasses and may offer a synthetic  
265 platform to mimic the oceanic carbon cycle. For practical field operation, the JPV can be scaled  
266 facilely and serve as a phototactic vehicle carrying other functional parts and electronic devices if  
267 necessary. For example, we demonstrate self-supported cycling in a series of  
268 microalgae-incorporated JPVs (**Supplementary Fig. S34 and Video 12**), in which the microalgae  
269 serves as the photothermal converter and can provide many biological and ecological functions.  
270 The animate-material system is also capable of sustainable cycling in a swarm, churning the  
271 ambient fluid under the cue of a constant spotlight (**Supplementary Fig. S35 and Video 13**).  
272  
273

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284

285 **Author contributions:** X.Q. conceived the concept, designed the experiment, and wrote the  
286 manuscript. X.Q., G.H. Z.L., F.D., W.L. carried out the material synthesis, characterization and  
287 systematic demonstration. G.H., Y.W., G.Y, X.Q., N.X.F. conducted the fluid dynamic experiment and  
288 analysis. G.H., X.Z, H.W., X.Q., carried out the numerical simulation. J.C, and G.M. supervised the  
289 device modeling. X.Q., and N.X.F supervised the project. All authors analysed and interpreted the  
290 data.  
291

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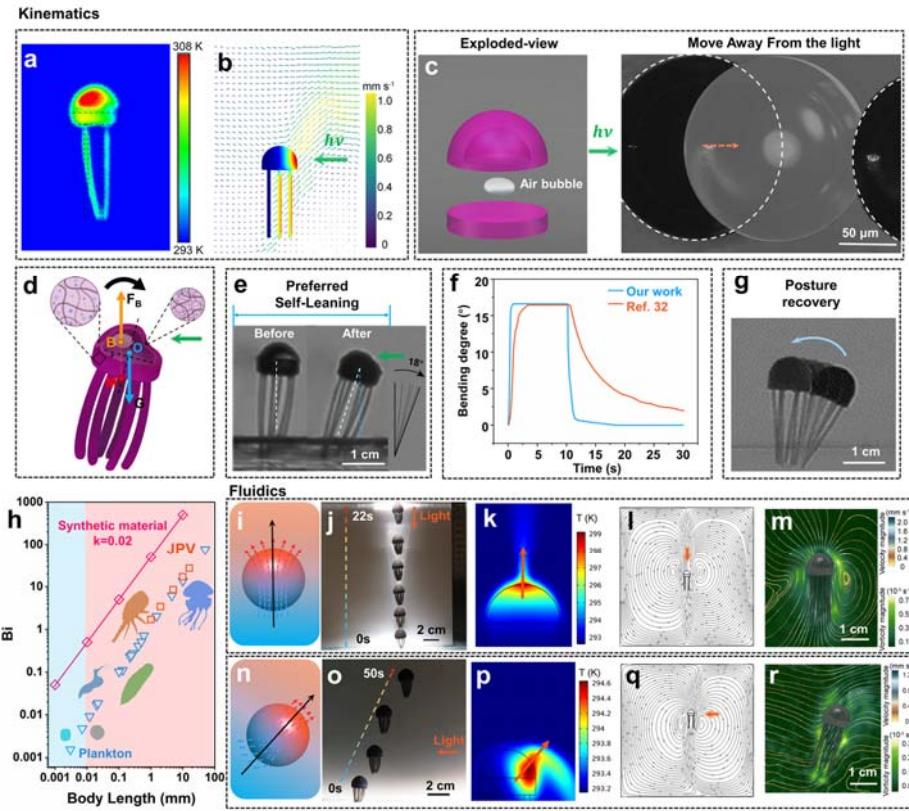
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301 **Fig. 1. Schematics of full-space phototaxis enabled by regulated symmetry-breaking.** a,  
 302 Simulated temperature gradient of the illuminated animate-material, due to the photo-thermal  
 303 coupling. The symmetry-breaking of temperature allows the material to sense, recognize and  
 304 move along the direction of photonic incidence.  $h\nu$  refers to the photonic incidence,  $\vec{v}$  the vector  
 305 of temperature gradient and moving direction. b, c, Schematics of the thermally induced fluidic  
 306 field (b) and thermo-mechanical coupling induced density redistribution (c), both are the results of  
 307 the symmetric breaking. d, Schematics of self-regulated locomotion that enables a negative  
 308 feedback loops and resulting e, Directional correction by the material itself. f, Schematics of the  
 309 fast recovery of the animate material system that is a key that enables the fast response to the  
 310 instantly changing direction of illumination. g, Schematic of the full-space phototaxis.

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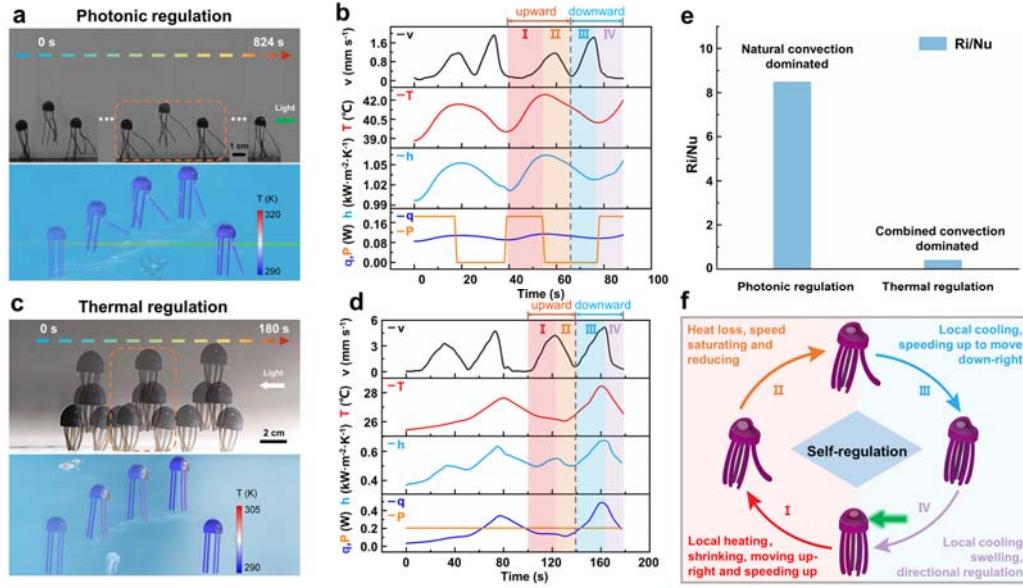
315 **Fig. 2. Photothermal-tactic flow-guided movement.** **a**, Infrared image of the JPV under  
 316 photonic irradiation. **b**, Snapshot of a PIV image showing the broken symmetry of the flow field,  
 317 in which the JPV model is included to highlight the location in the PIV experiment and the  
 318 temperature gradient. **c**, Schematic exploded-view and photographic image of the embedded  
 319 bubble of the JPV. The bubble in the bell moves away from the irradiation from the left. **d**, Local  
 320 shrinkage squeezes the bubble and deviates the center of force to produce overturning moment. **e**,  
 321 Self-leaning effect of  $18^\circ$  due to symmetric breaking induced by the relocation of the bubble. **f**,  
 322 Actuation and recovery rate (for  $20^\circ$  bending) of the ultrafast hydrogel designed for the JPV and  
 323 comparison with previously reported data. **g**, Posture recovery when the symmetries of the JPV  
 324 are restored. **h**, Comparison of Biot number of plankton, synthetic material and JPV. **i**, Schematic  
 325 illustration of vertical illumination, where the heating of the top brings about an upward fluid  
 326 motion. The arrows indicate the direction of the photothermal tactic flow. **j**, Composite image  
 327 showing the JPV ascending from the bottom to the water surface under vertical illumination. **k**, **l**,  
 328 FEM-simulated temperature gradient (**k**) and flow field (**l**) around the JPV under vertical  
 329 illumination. **m**, Magnitude of velocity and vorticity recorded by PIV flow measurement during  
 330 vertical illumination. **n**, Schematic illustration of asymmetric illumination, where local heating  
 331 generates asymmetric fluid motion. The arrows indicate the direction of photothermal tactic flow.  
 332 **o**, Composite image showing the movement of the JPV under horizontal illumination. **p**, **q**,  
 333 FEM-simulated temperature gradient (**p**) and flow field (**q**) around the JPV under horizontal  
 334 illumination. **r**, Magnitude of velocity and vorticity recorded by PIV flow measurement during  
 335 horizontal illumination. Input photonic power density:  $2 \text{ mW} \cdot \text{mm}^{-2}$ .

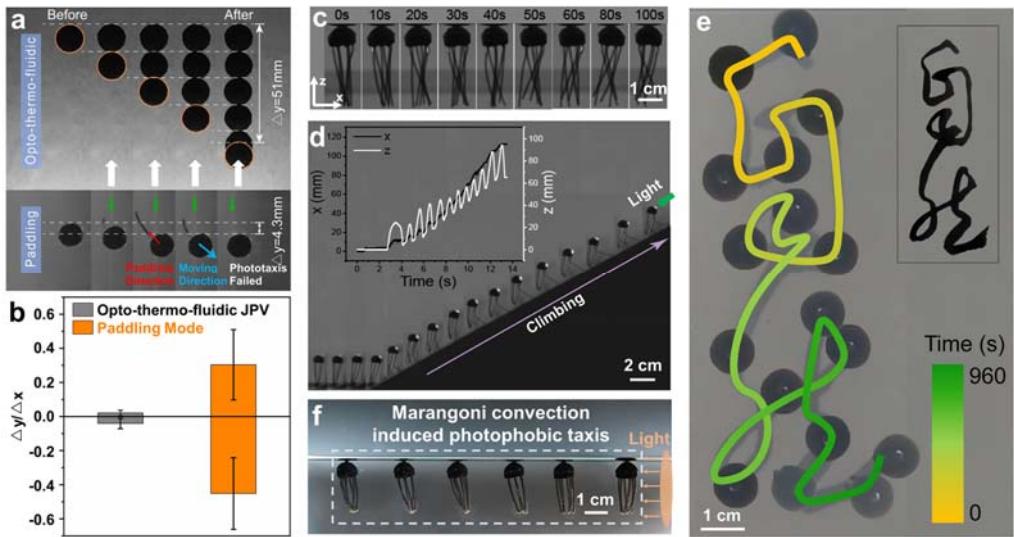
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337 **Fig. 3. Phototaxis of the JPV enabled by self-regulated symmetric breaking.** **a**, Composite  
 338 snapshots and FEM analysis of phototactic hopping towards a laser source (532 nm, 200 mW) **b**,  
 339 Time evolutions of scalar velocity, temperature, convective heat transfer coefficient, input  
 340 photonic power and convective cooling power near the illuminated bell area of the JPV over two  
 341 consecutive cycles. **c**, Composite snapshots and FEM analysis of phototactic hopping of the JPV  
 342 towards a spotlight (beam diameter of 10 cm, 2 mW  $\cdot$  mm $^{-2}$ ). **d**, Time evolutions of scalar velocity,  
 343 temperature, convective heat transfer coefficient, input photonic power and convective cooling  
 344 power near the illuminated bell area of the JPV over consecutive cycles. **e**, Comparison of  
 345 Richardson number and Nusselt number for photonic and thermal regulation when velocity is at its  
 346 peak. **f**, Summarized mechanism of the phototactic locomotion. The JPV's motion automatically  
 347 introduces negative feedback loops that regulate photonic and thermal-energy-induced symmetric  
 348 breaking.

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353 **Fig. 4. Accurate and sophisticated motions of underwater phototactic steering.** **a**, Comparison  
 354 of the phototaxis accuracy in the photo-thermal-fluidic mode (top) and tentacle-paddling mode  
 355 (bottom). The former achieves travel over at least 51 mm (limited by the size of the container),  
 356 whereas the latter exhibits a failure to move after 4.3 mm due to misalignment between the  
 357 tentacle and the laser beam. **b**, Analysis of phototactic accuracy by quantifying the deviated  
 358 displacement along the y-axis over the travelling distance in the x direction ( $\Delta y/\Delta x$ ). **c**,  
 359 Snapshots showing the rotational advancing of the JPV by cascading the strokes of different  
 360 tentacles. **d**, Composite snapshots and trajectory curve (inset) depicting the phototactic climbing  
 361 of the JPV in x and z directions. **e**, JPV manœuvred with finesse, accurately representing each  
 362 stroke of the Chinese calligraphic characters “自然”. **f**, At the water surface, the JPV shies away  
 363 from the light due to Marangoni convection.

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